

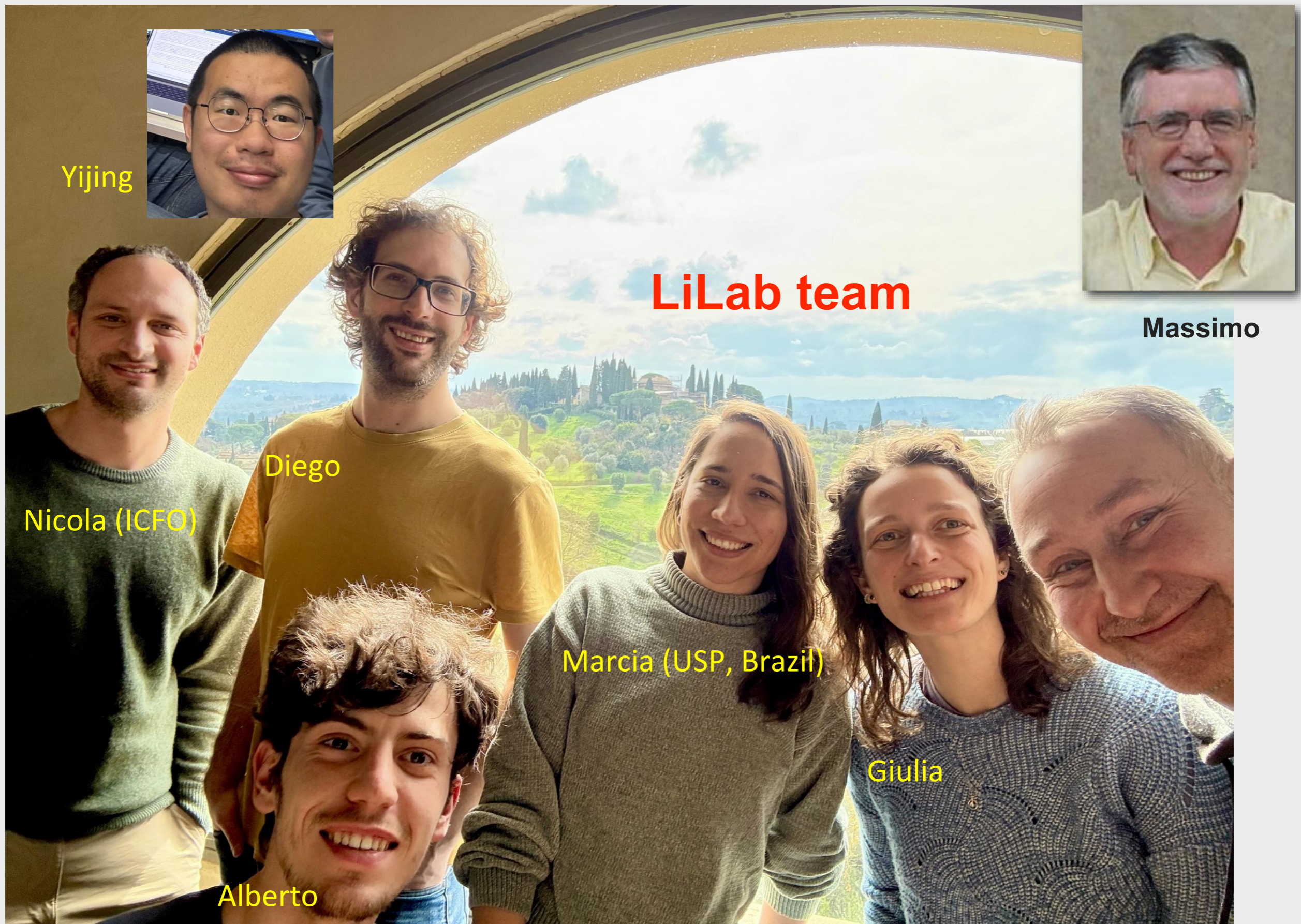
Tunnelling transport in strongly-interacting atomic Fermi gases

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June 18th 2026

ICAP2026, Wuhan, China





Yijing



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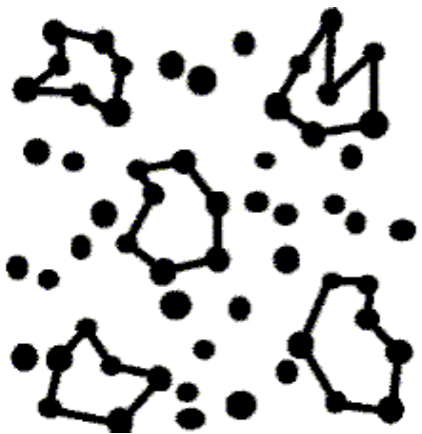
Why quantum transport?

Transport measurements: fundamental tool to investigate matter.

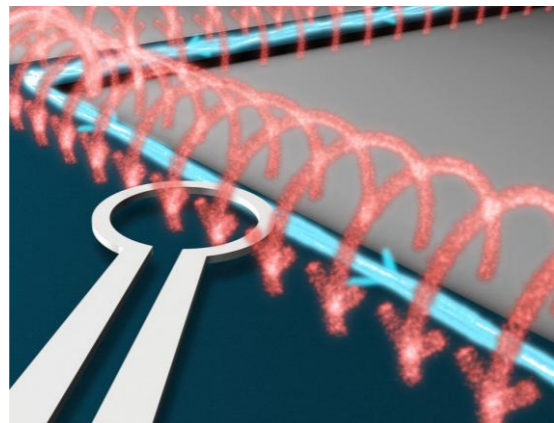
Extracting microscopic properties of a quantum many-body systems via global measurements, typically current (T. Giamarchi).

Datta, *Electronic transport in mesoscopic systems*
(Cambridge University Press, 1997)

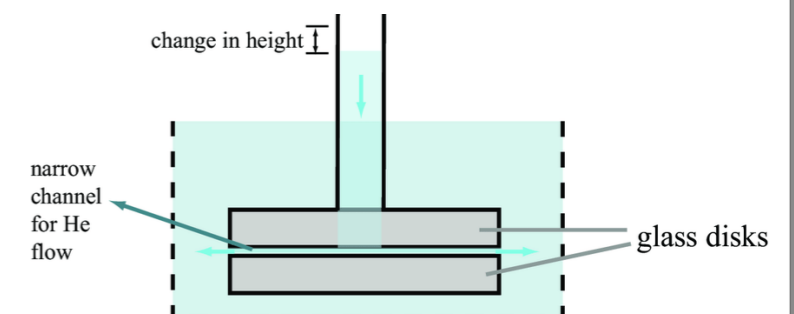
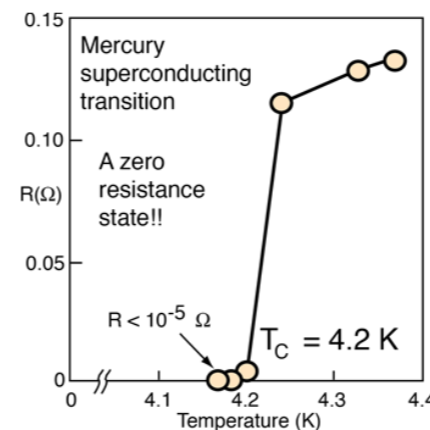
Localization



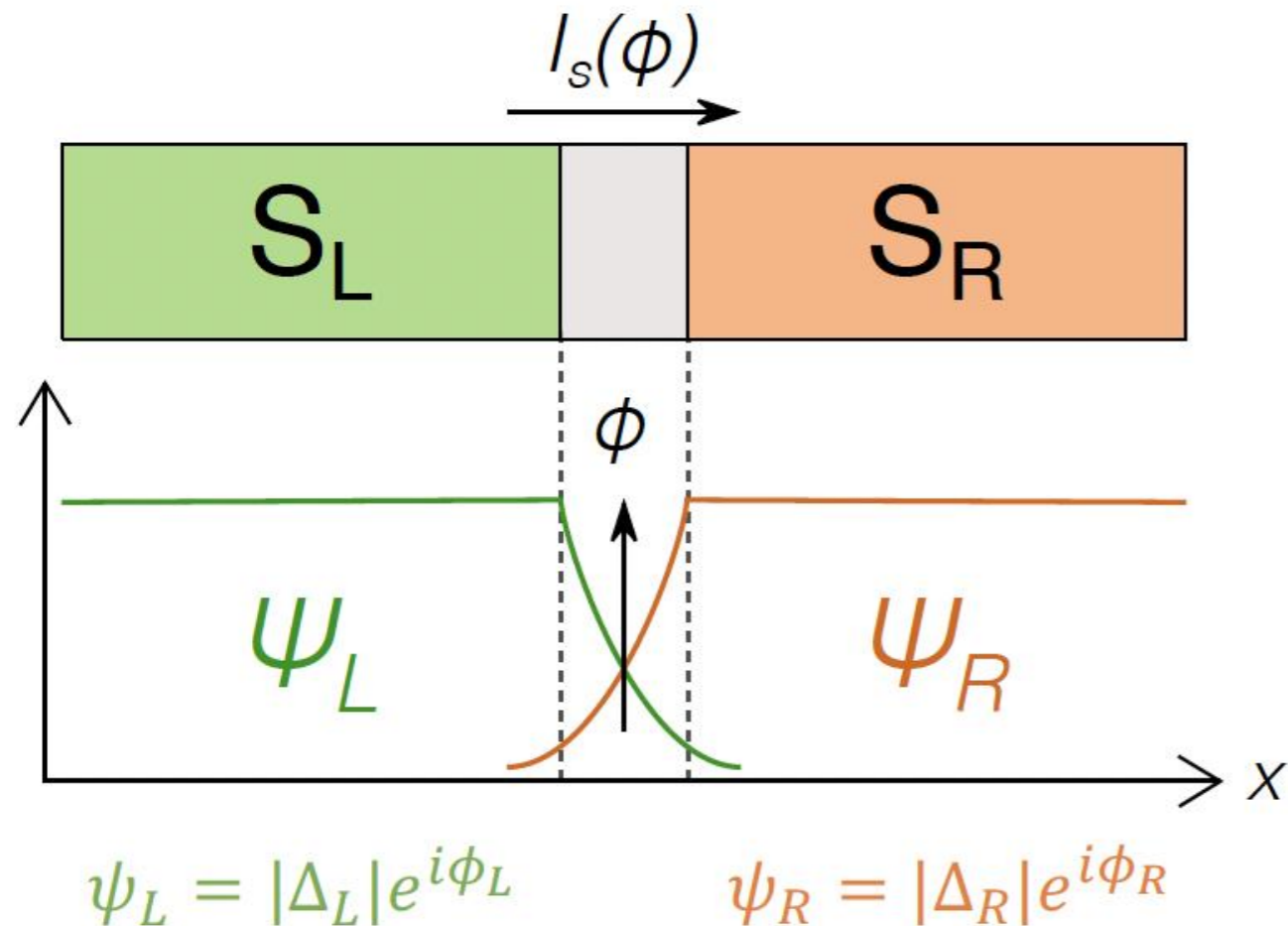
Topological currents



Superconductivity and superfluidity



A paradigmatic example: a Josephson junction



Josephson-Anderson equations:

$$I = I_c \sin(\phi)$$
$$\dot{\phi} = -\frac{2e}{\hbar} V$$

$\phi = \phi_L - \phi_R \rightarrow$ relative phase

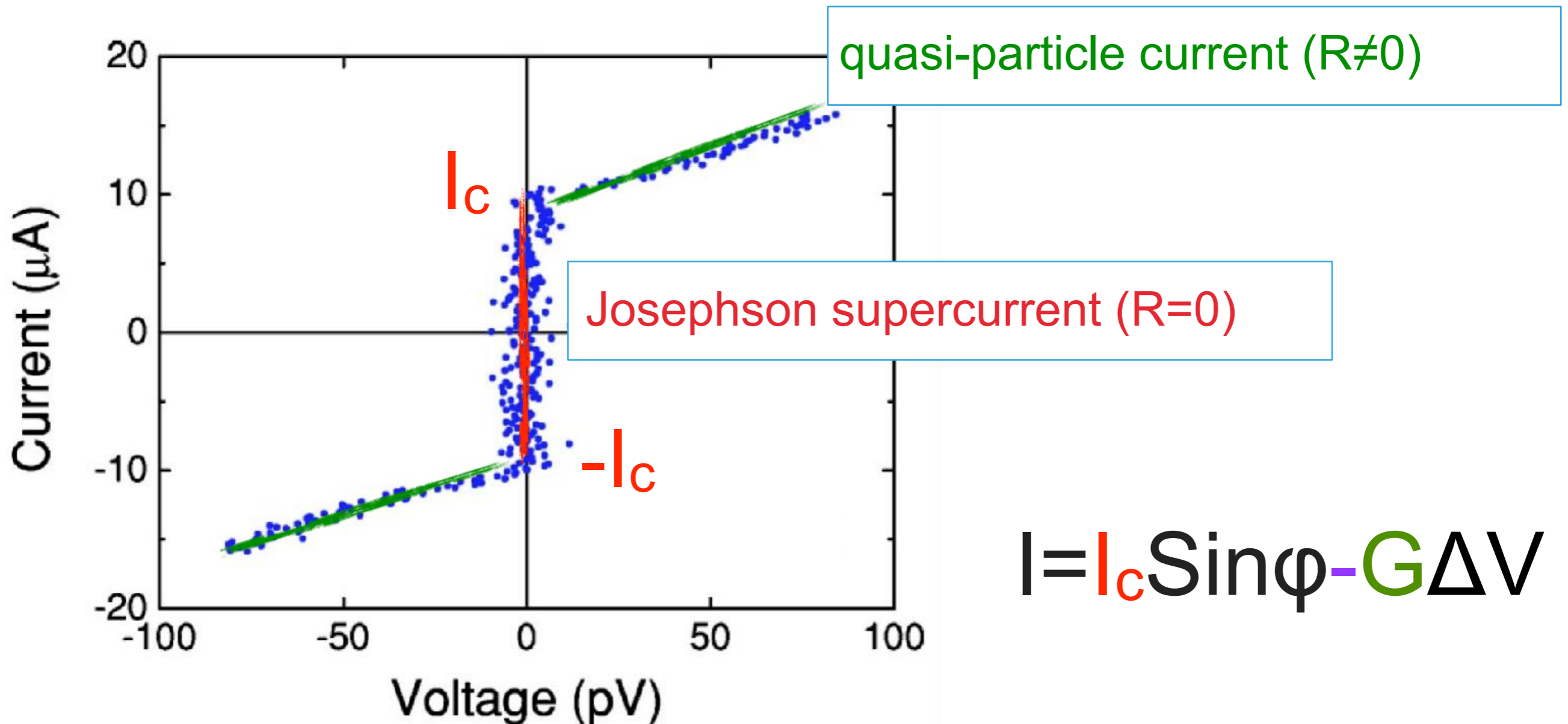
$I_c \rightarrow$ critical current

$V \rightarrow$ potential difference

☑ Thin barrier: quantum tunneling regime

$\varphi = \varphi_1 - \varphi_2$: “quantum voltage”

$$I = I_c \sin \varphi$$



$$I = I_c \sin \varphi - G \Delta V$$

Why Josephson tunneling?

Ambegaokar-Baratoff relation (BCS)

BEC-BCS crossover: Zaccanti and Zwerger PRA (2019)

$$\hbar I_c \equiv E_J \sim \Delta \sim n c$$

Measurable current (macro)



Condensate fraction (micro)

(challenging for strongly-interacting gases)

The Josephson effect (Anderson's viewpoint)

... In the case of superconductivity things are quite different. The internal long-range order parameter — the phase — is not a parameter for which suitable instruments exist; we do not normally walk around with objects which have phase order. A superconductor has rather perfect internal phase order, but as we have shown the zero-point motion of the total order parameter of an isolated superconductor is large and rather rapid.

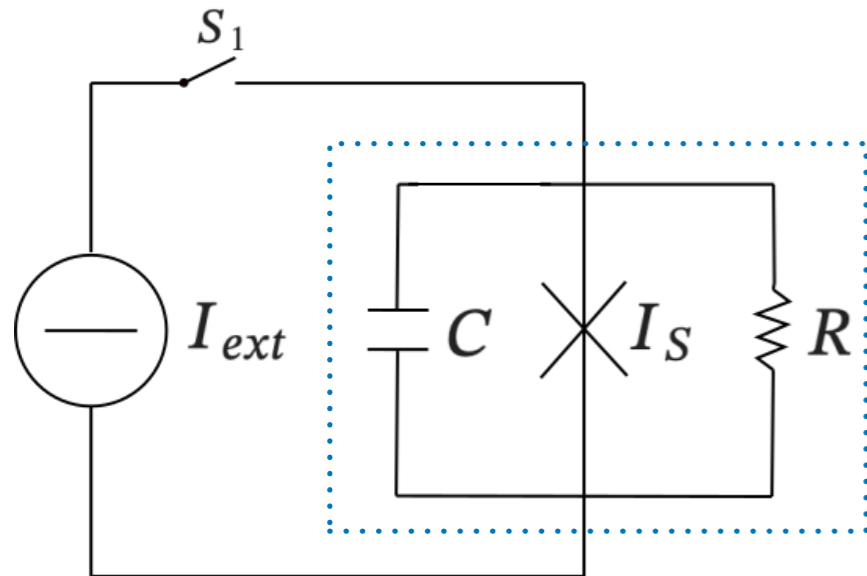
The importance of the Josephson effect, then, is that it provides for the first time an instrument which can act like a clamp for a solid or a coercive field for a ferromagnet: it can pin down the order parameter ...

But most important of all, it frees us conceptually from the mystery surrounding superconductivity, and places it in line with all of the other condensation phenomena.

Anderson in Lectures on the Many-body Problems, **E.R. Caianiello, Elsevier Science (1964)**

RCSJ model: an operative mechanical model

Tinkham, Introduction to Superconductivity



Kirchhoff law:

$$I_{ext} = I_c \sin(\phi) + GV + C\dot{V} + \dot{\phi} = \frac{2e}{\hbar} V$$

$$\frac{C\hbar}{2e} \ddot{\phi} = I_{ext} - I_c \sin(\phi) - \frac{G\hbar}{2e} \dot{\phi}$$

$C \rightarrow$ capacitance

$R \rightarrow$ resistance

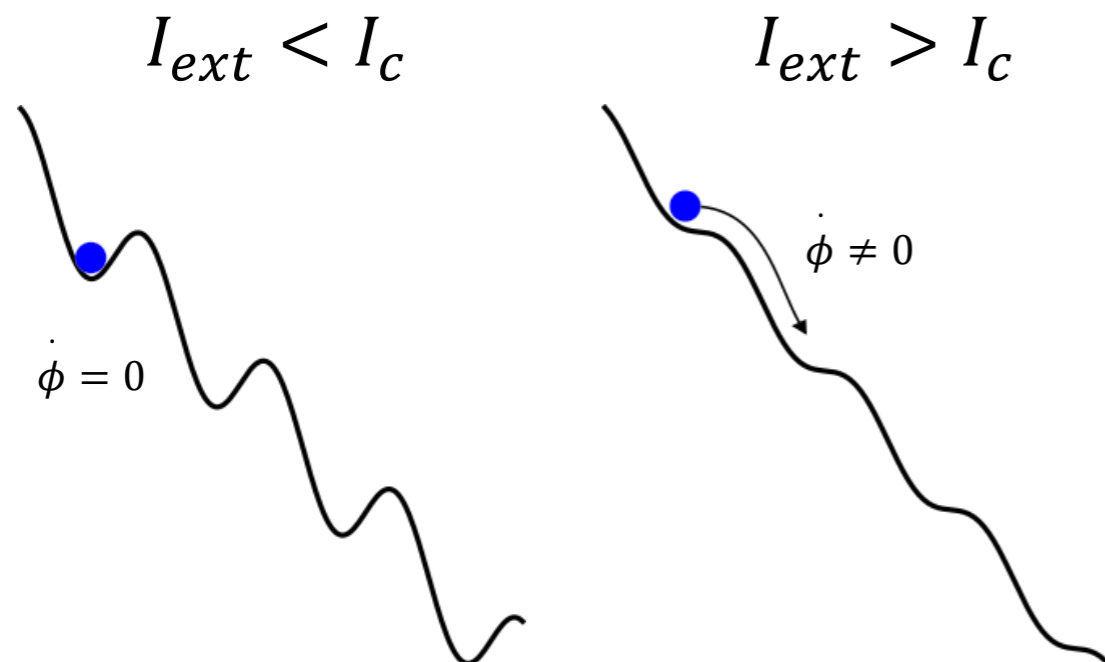
$G = 1/R \rightarrow$ conductance

Equation of motion for a phase particle of mass $\sim C$ subject to a viscous force $\sim G$ in the **washboard potential**:

$$U(\phi) = -I_{ext}\phi + I_c \cos(\phi)$$

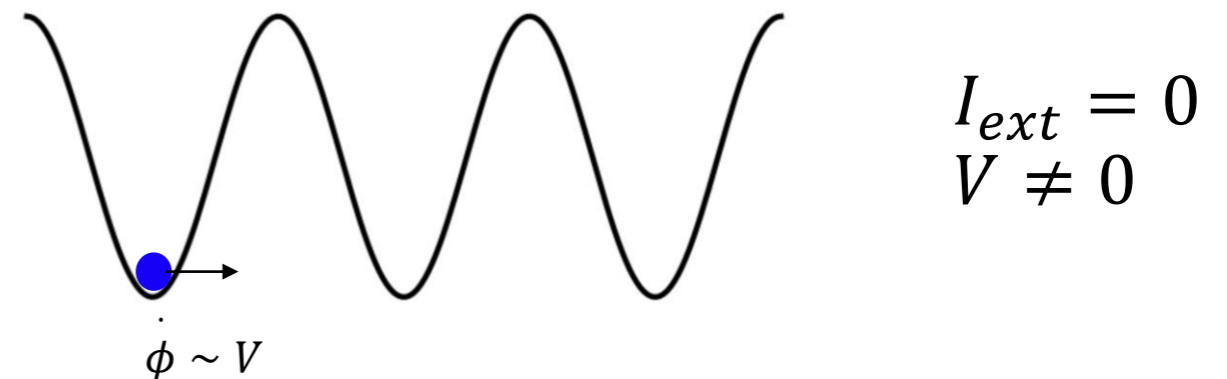
Washboard potential

DC Josephson effect



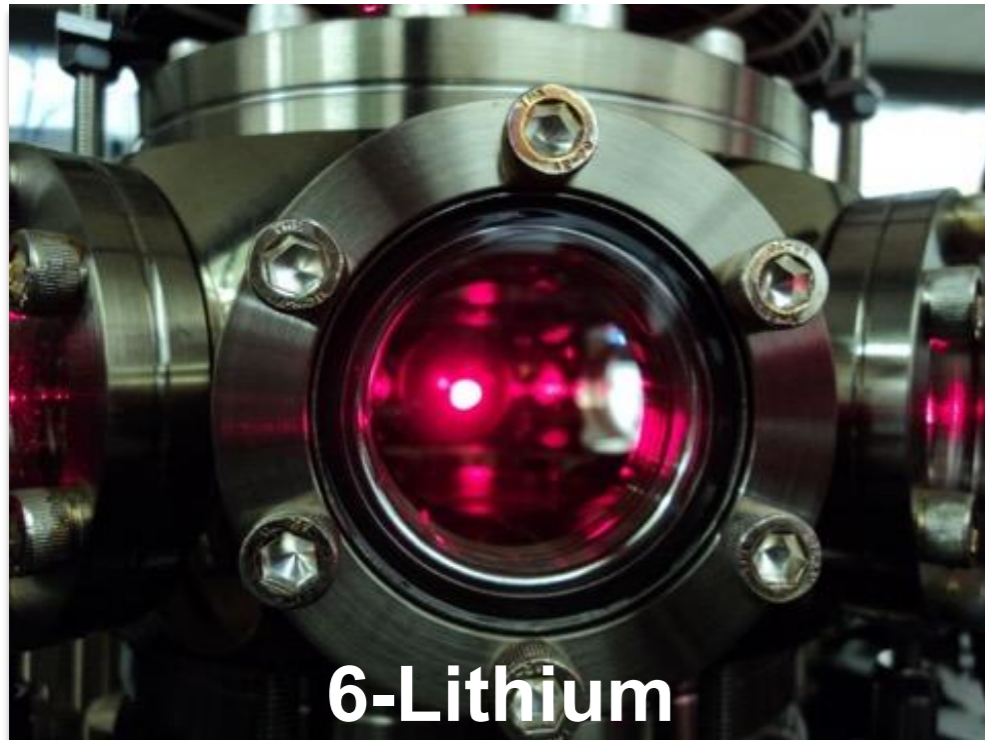
$I_{ext} < I_c$: the phase particle remains in one washboard potential minima: dissipationless current.

AC Josephson effect

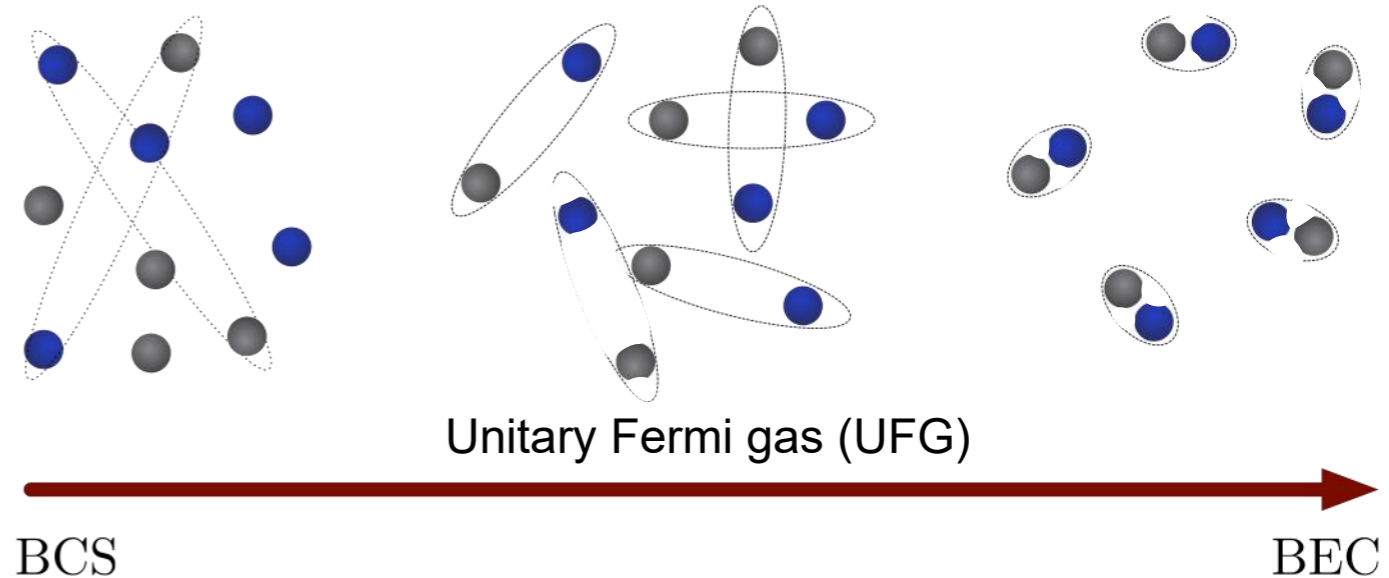
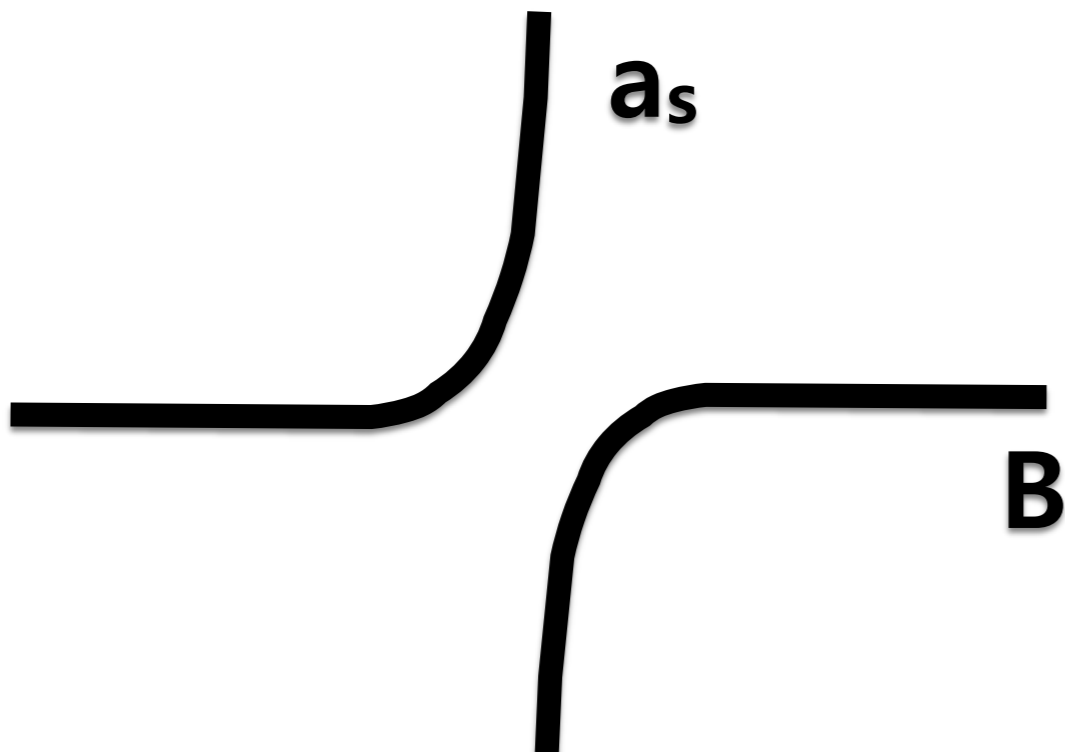
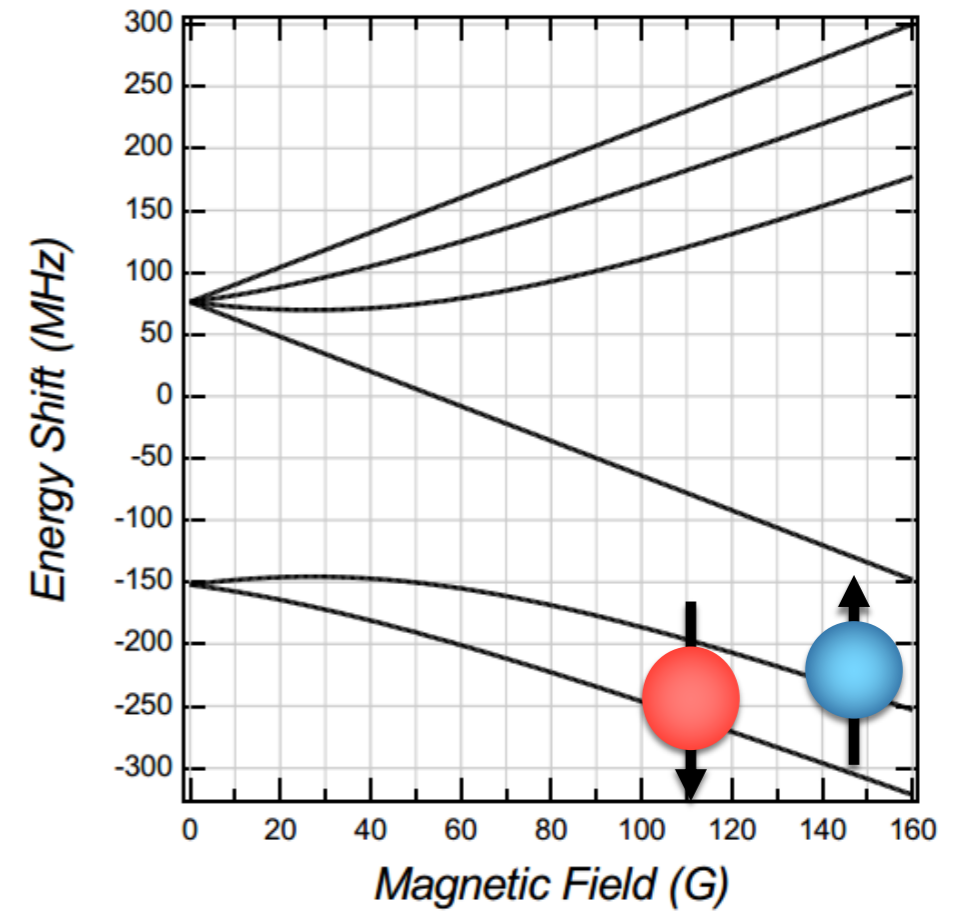


The phase particle oscillates in one washboard potential minima: Josephson oscillations @ $\omega_J \sim \Delta V / \hbar$

^6Li and BEC-BCS crossover



+

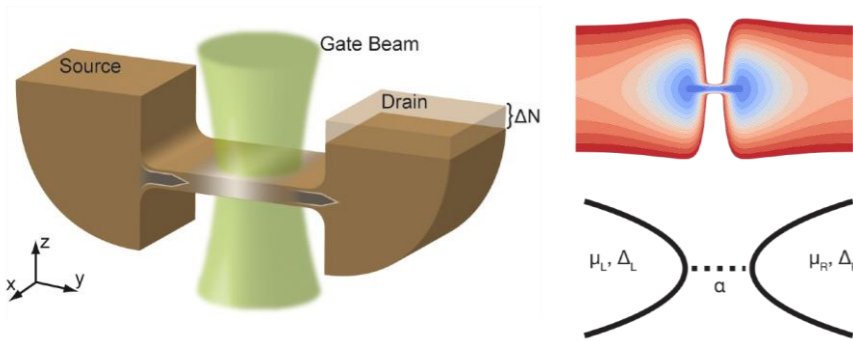


Holland, Kokkelmans, Chiofalo, and Walser PRL (2001)

Randeria, Nat. Phys. (2010)

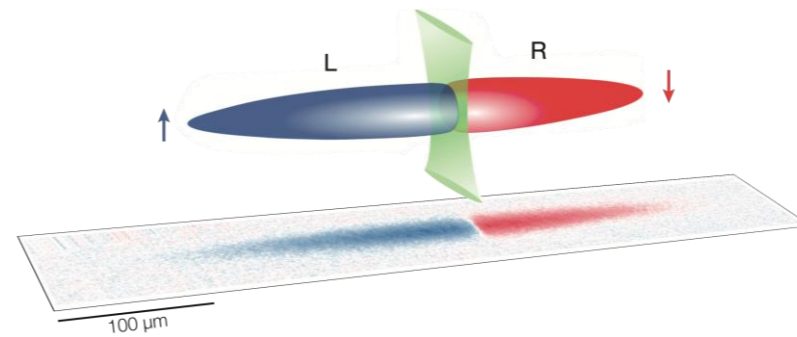
Quantum Transport with Fermi Gases

Transport through narrow ballistic channels



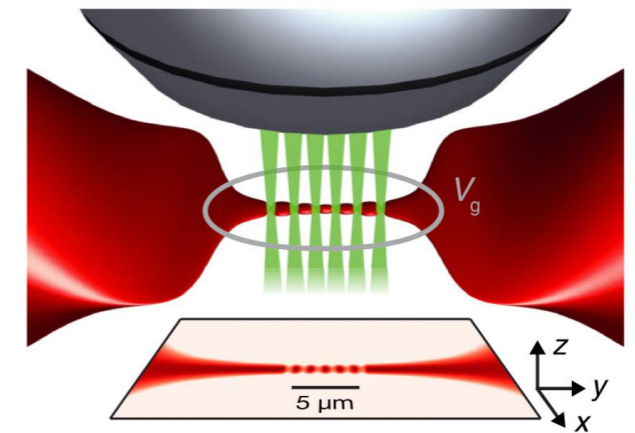
ETH, Geneva

Spin, sound and heat transport..



MIT, Harvard, Innsbruck, Toronto,
ETH, LENS, MPQ, Rice,
Hamburg,...

Tunneling through structures



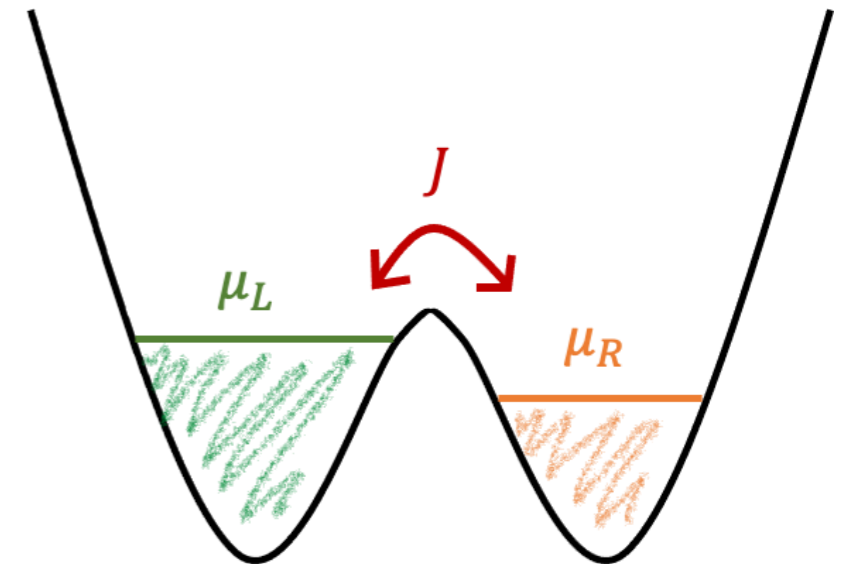
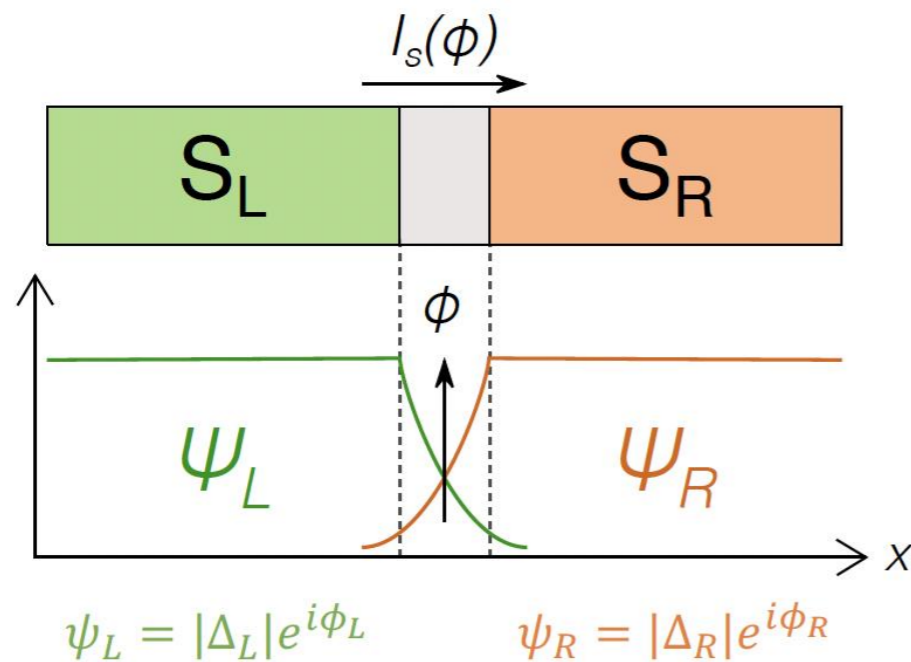
ETH,
Toronto,...

LENS,

Princeton,

For a review: [Krinner et al., J. Phys.: Condens. Matter](#) **29** (2017)

Electrons vs neutral atoms



Josephson-Anderson equations:

$$I = I_c \sin(\phi)$$

$$\dot{\phi} = -\frac{2e}{\hbar} V$$

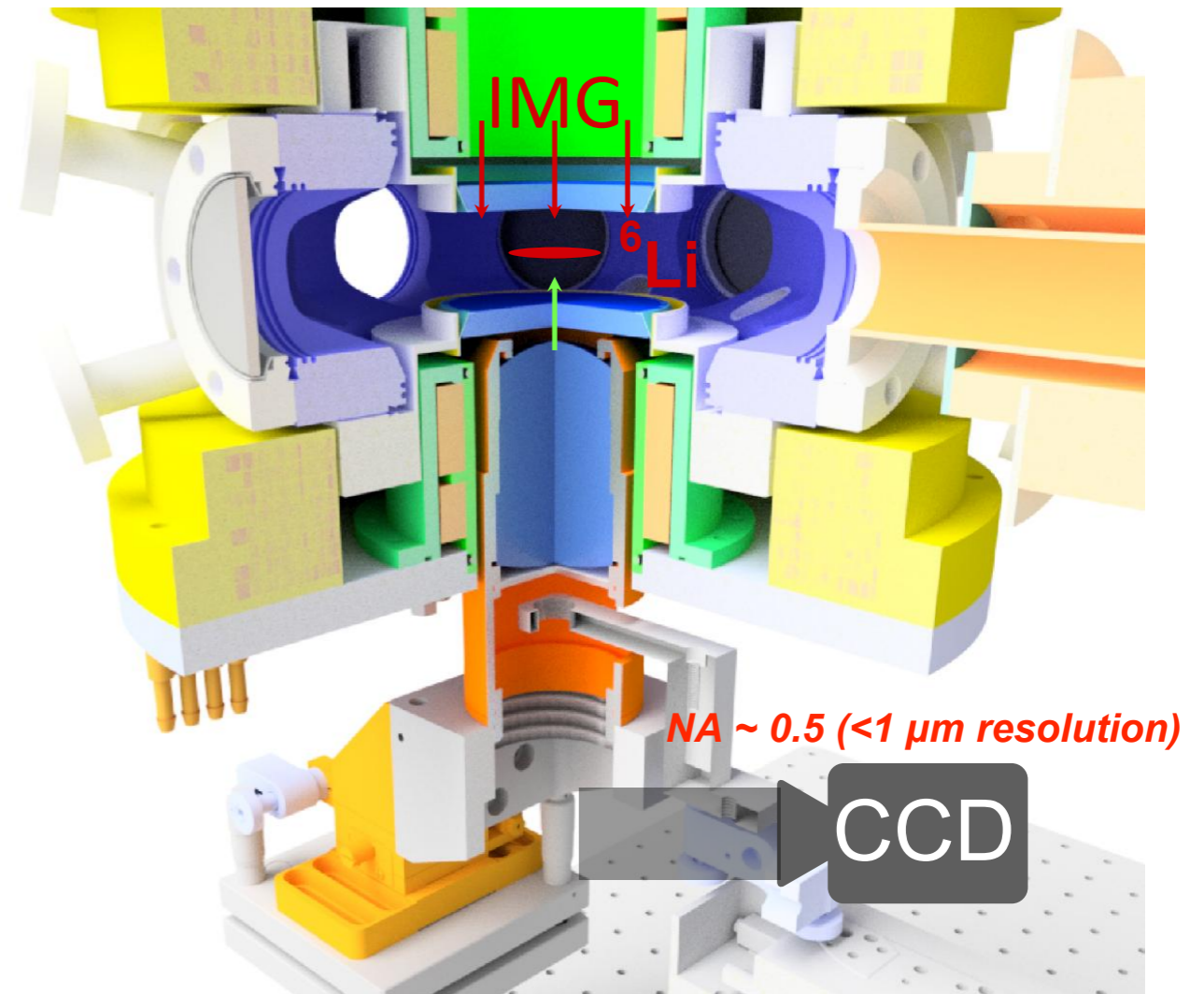
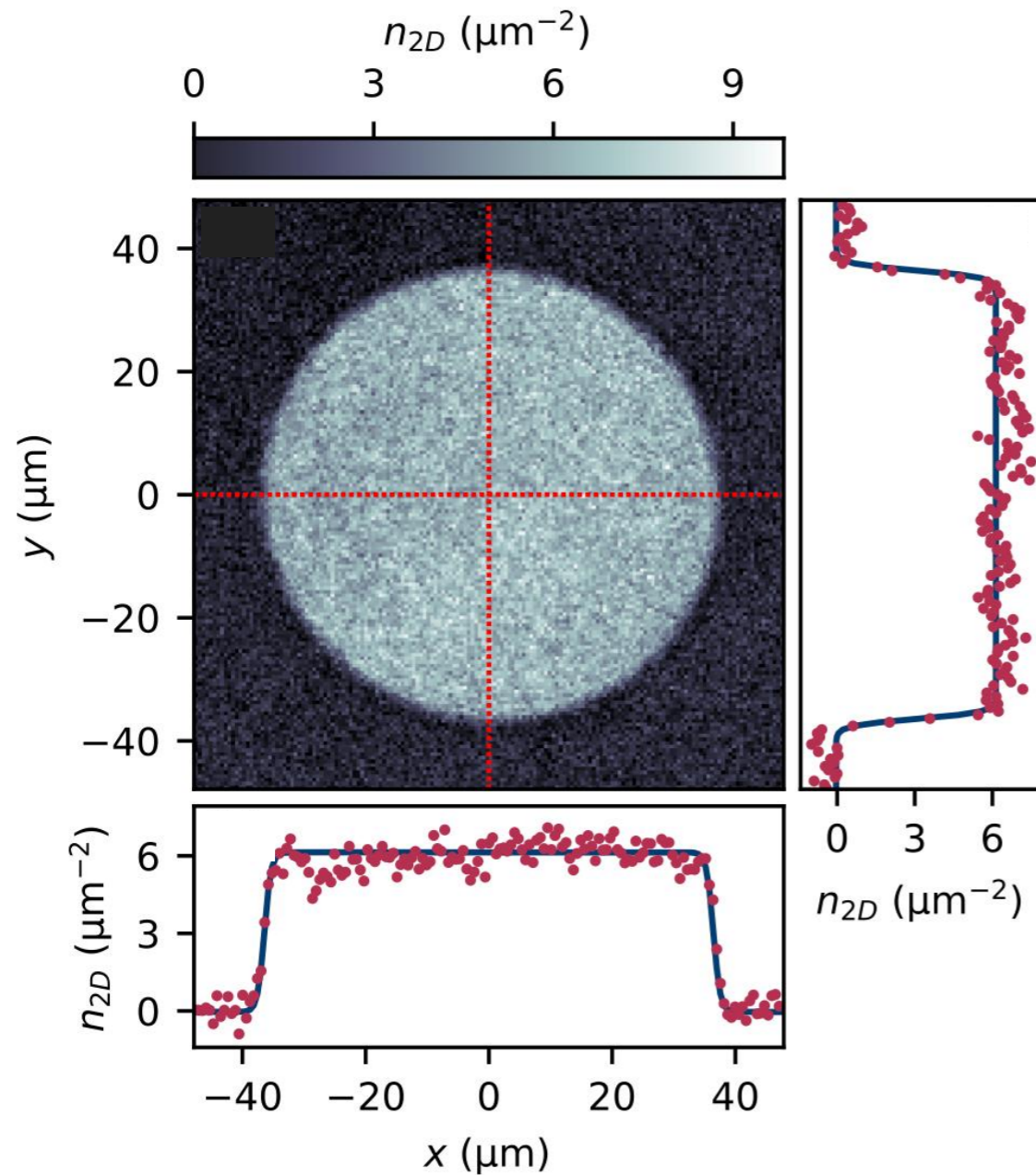
$$\phi = \phi_L - \phi_R \rightarrow \text{relative phase}$$

$$I = I_c \sin(\phi)$$

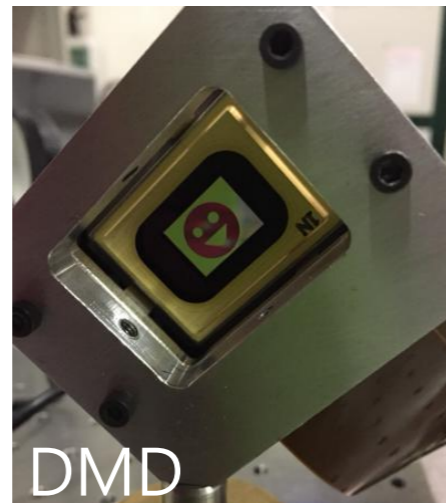
$$\dot{\phi} = -\frac{\Delta\mu}{\hbar}$$

Atomic Fermi gases @ LENS

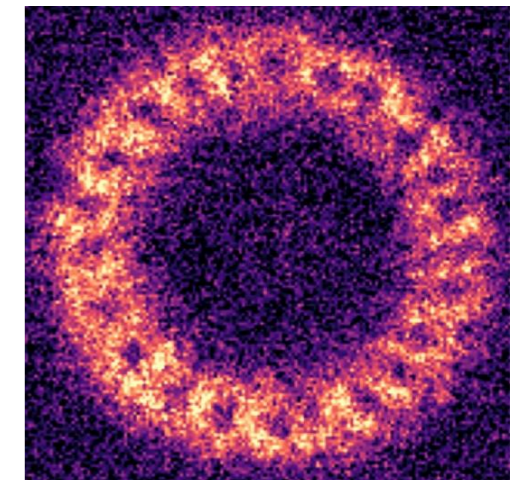
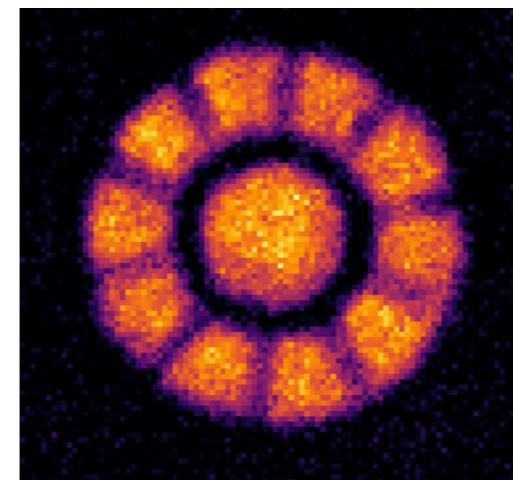
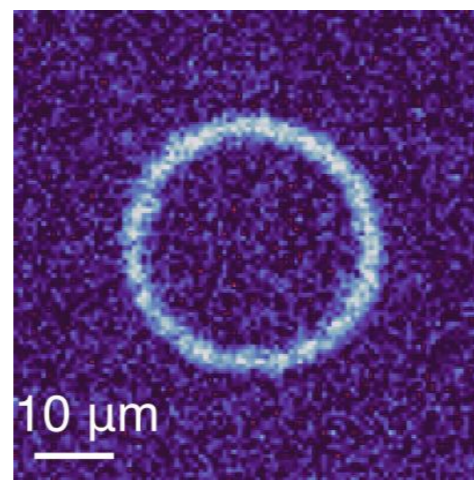
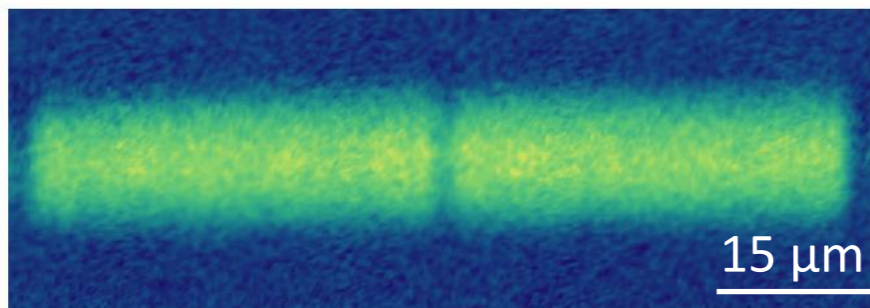
Quasi-homogeneous 3D superfluids: $N_{\uparrow}=N_{\downarrow}\approx 50000$ @ $T/T_F\sim .1 \rightarrow T < T_C$



Ingredients: shaping dynamical optical potentials



=

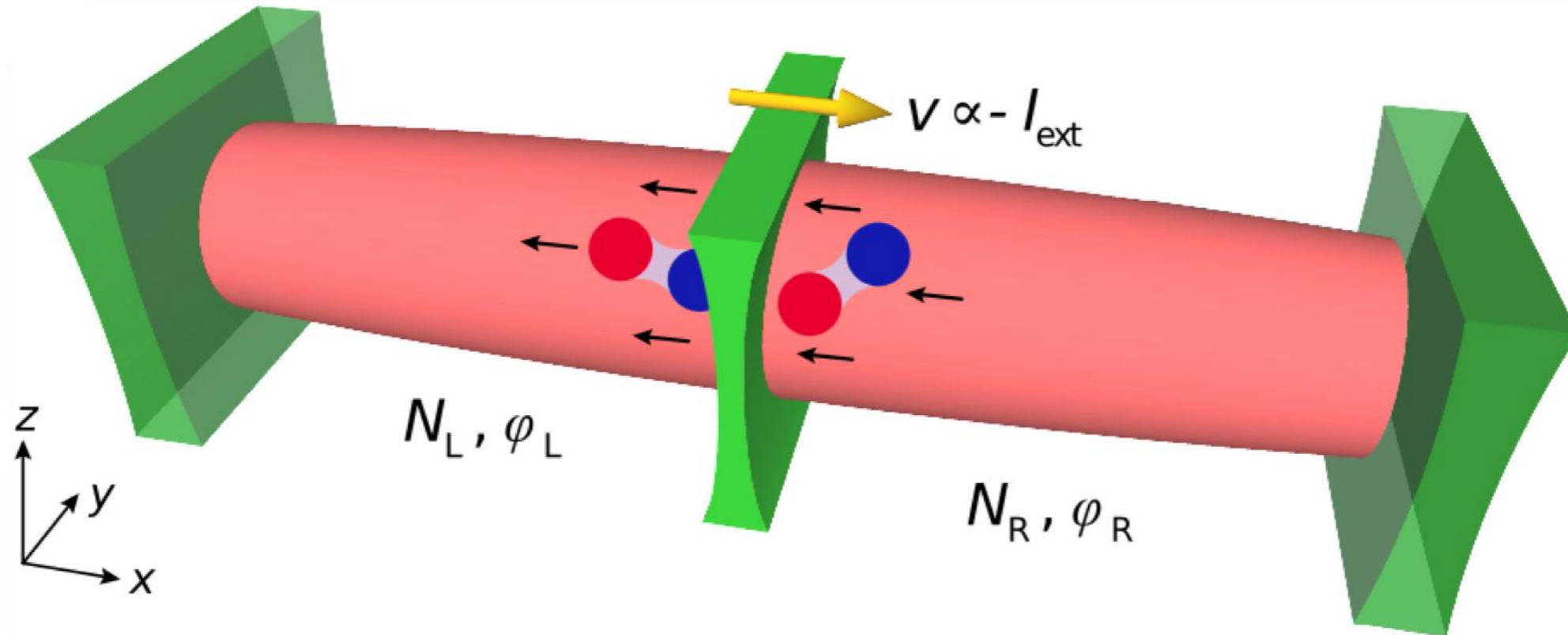


Valtolina et al., Science (2015)
Kwon et al., Science (2020)
Del Pace et al., PRL (2021)
Del Pace et al., Science (2025)

Del Pace et al., PRX (2022)
Pezze' et al., Nat. Comm. (2024)
Frometa et al., arXiv:2511.02664v2 (2025)
Ciszak, et al., arXiv:2601.15121v2 (2026)

Kwon et al., Nature (2021)
Rajkov et al., Nat. Phys. (2024)
Grani et al., Nat Comm. (2025)
Grani et al., Eur. Phys. J. (2025)

Injecting a current: moving the optical barrier at controllable velocity



Atomic cloud:

$$N_{R,L} \simeq 3.5 \times 10^4 \text{ atom pairs}$$

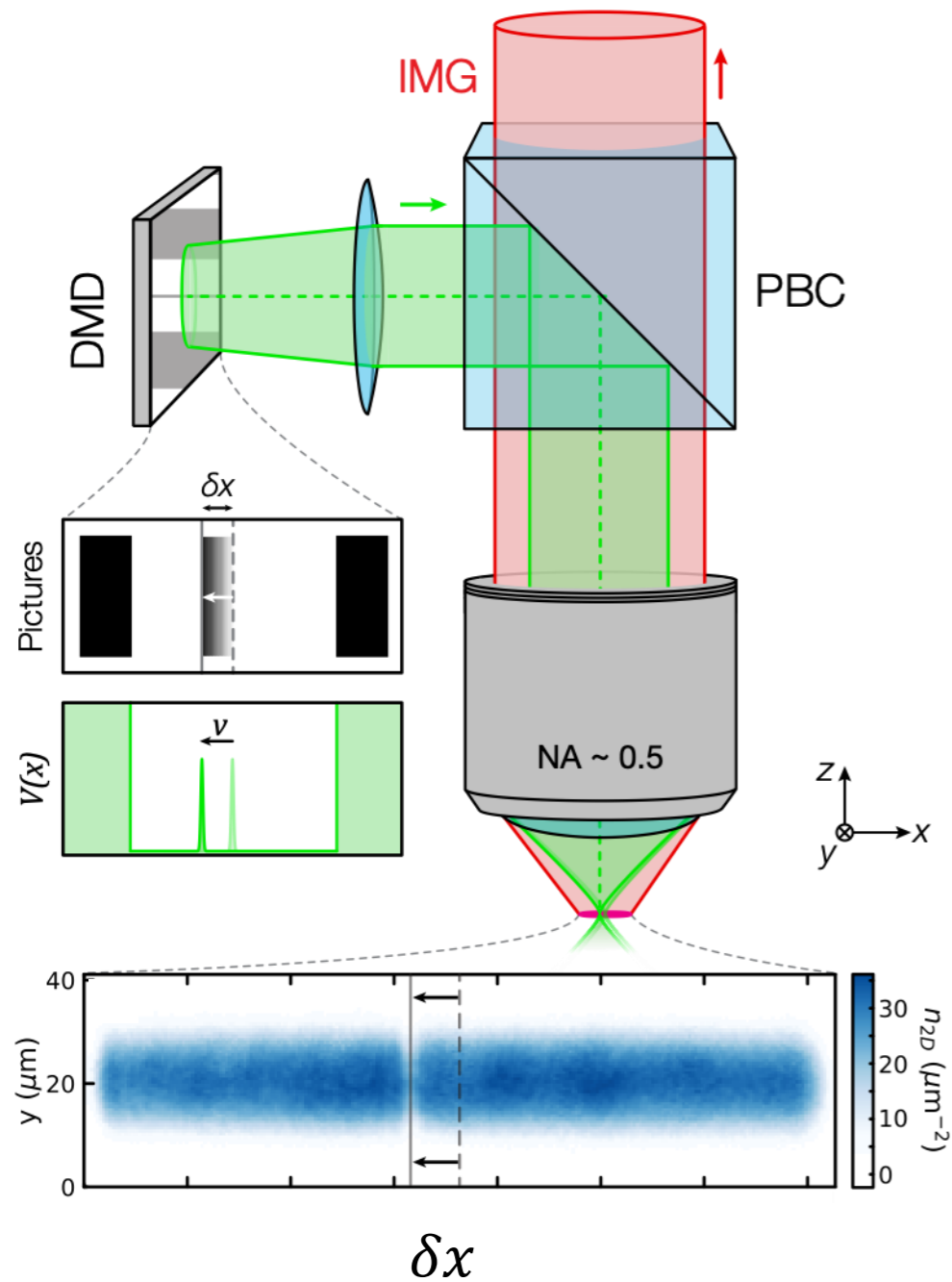
$$T/T_F \simeq 0.06 \text{ (2)}$$

Tunneling barrier:

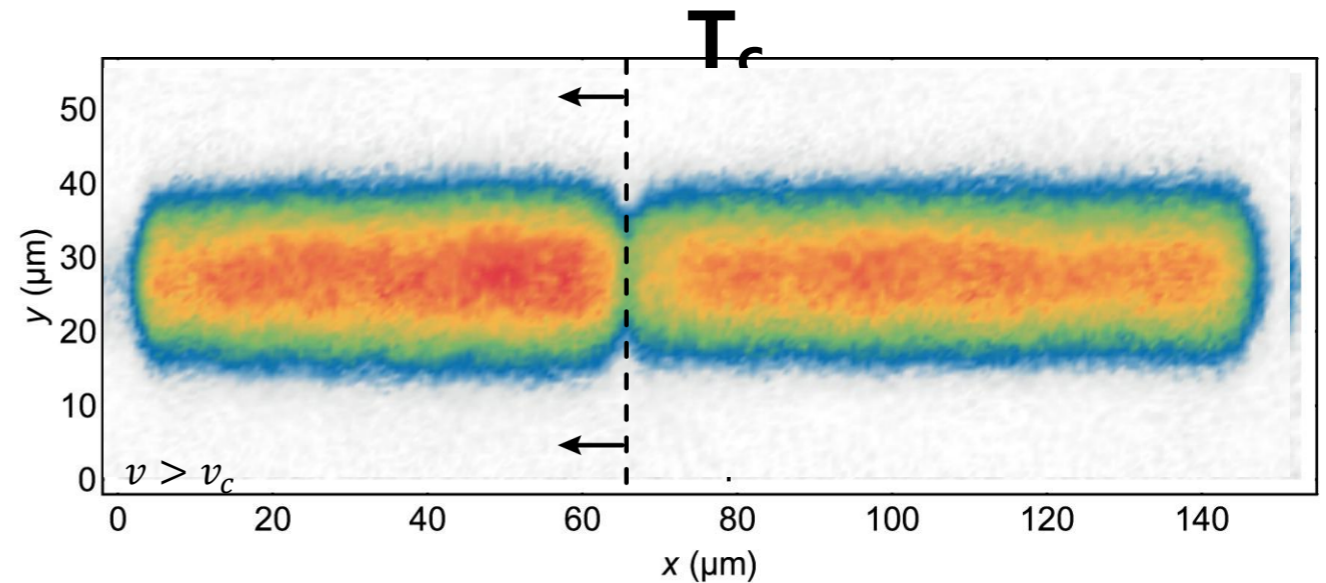
532 nm light

$$w \simeq 0.95 \text{ (10)} \mu\text{m}$$

$$V_0 \gtrsim \mu$$



3D unitary gas at $T \sim 0.3$



Below a critical current, the gas flows entirely through the tunnelling barrier → Superfluid coherent transport!

$$\delta x = 10 \mu\text{m}$$

$$v = 0.1 \text{ mm/s}$$

$$\sigma = 1.5 \mu\text{m}$$

$$V_0/\varepsilon_F \approx 0.8$$

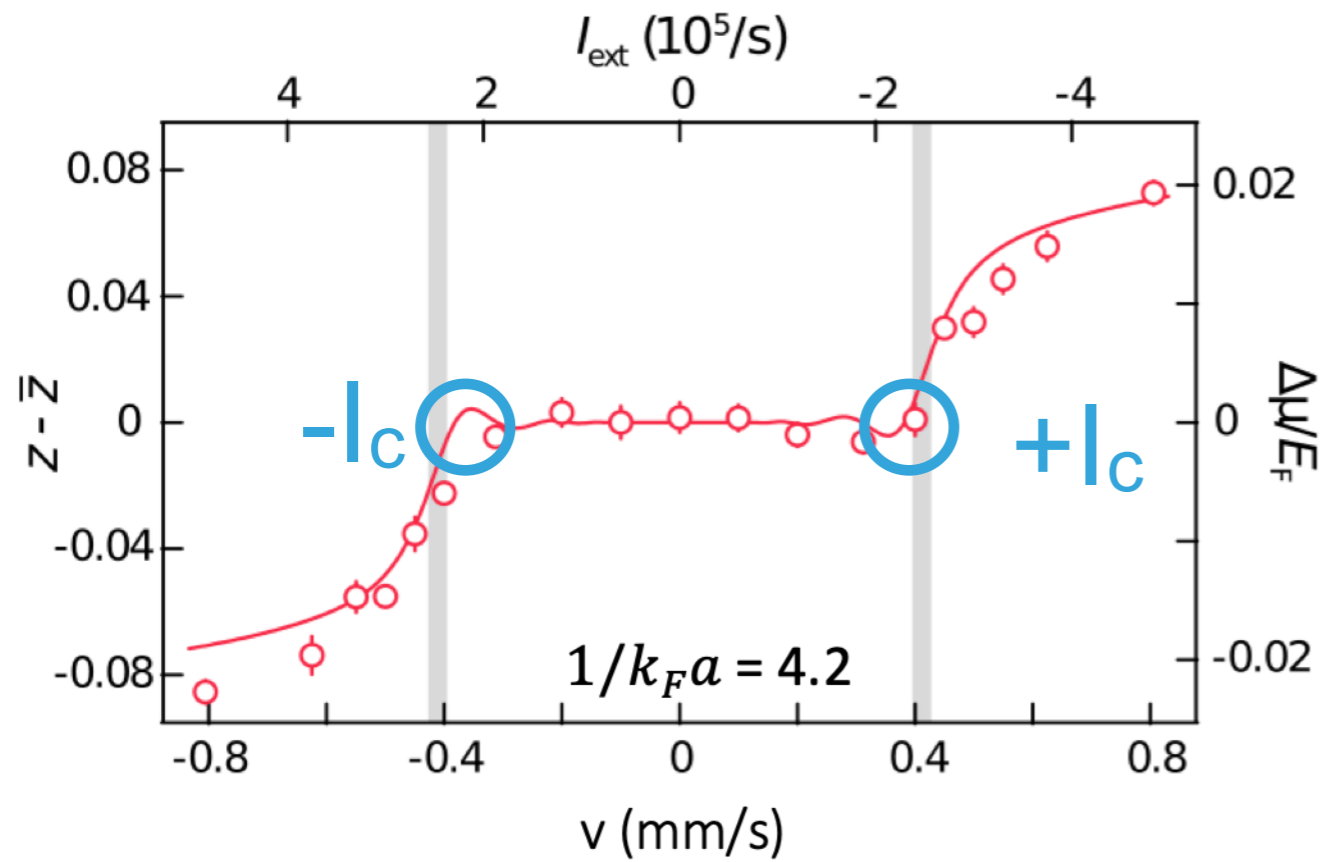
Sound velocity (bulk):
 $c \approx 10 \text{ mm/s}$



RCSJ circuit model

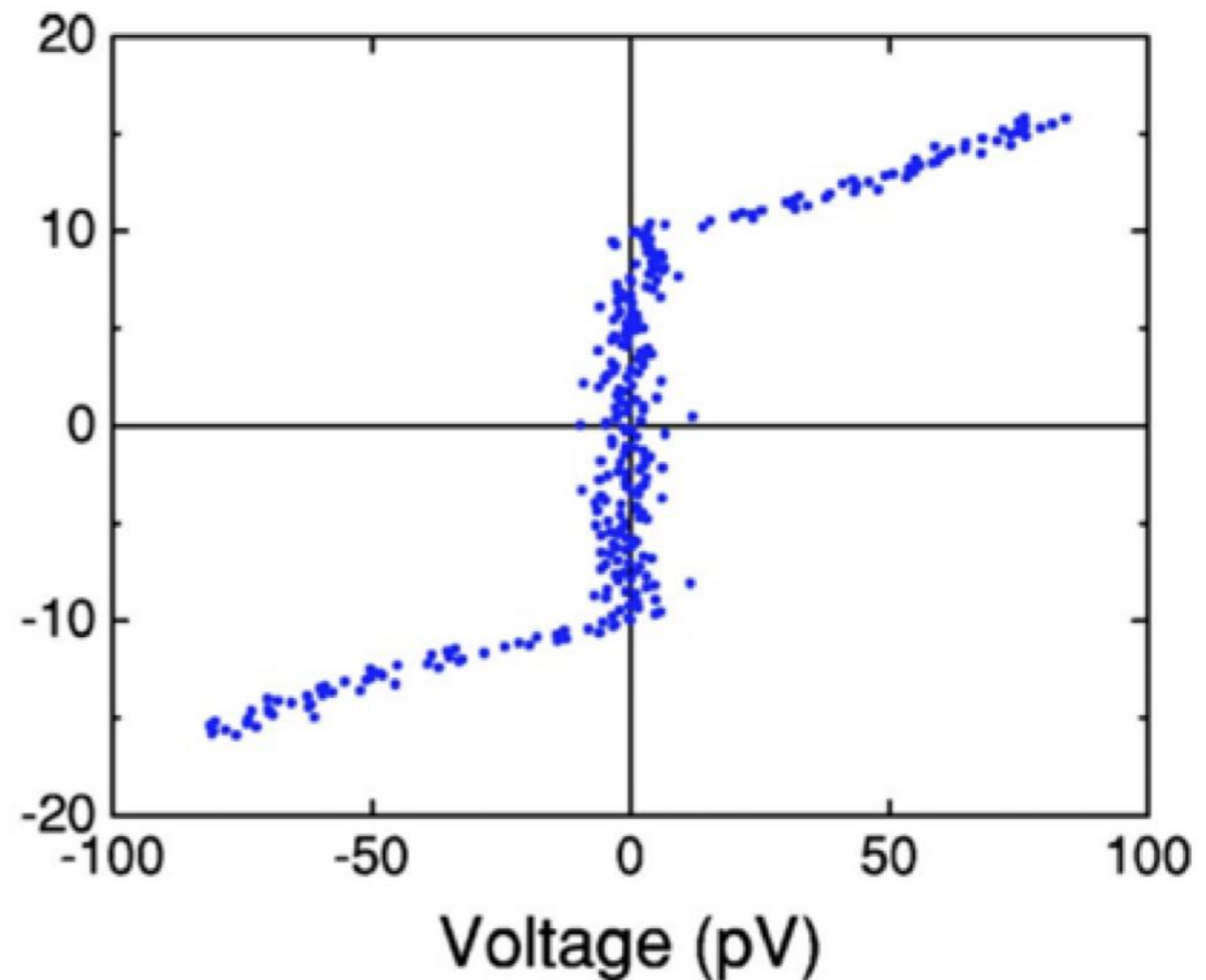
M. Tinkham, Introduction to Superconductivity

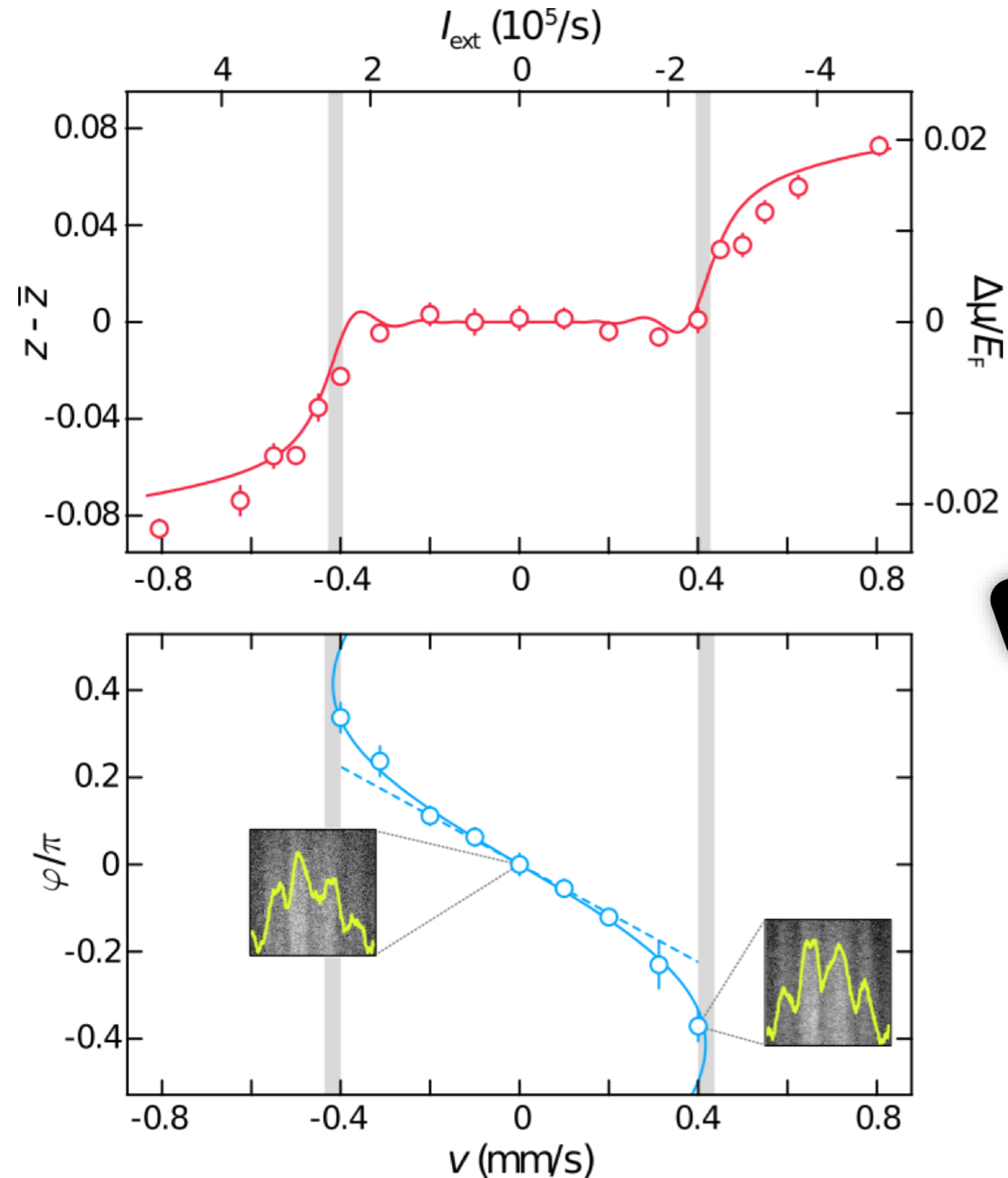
I_c and G free parameters

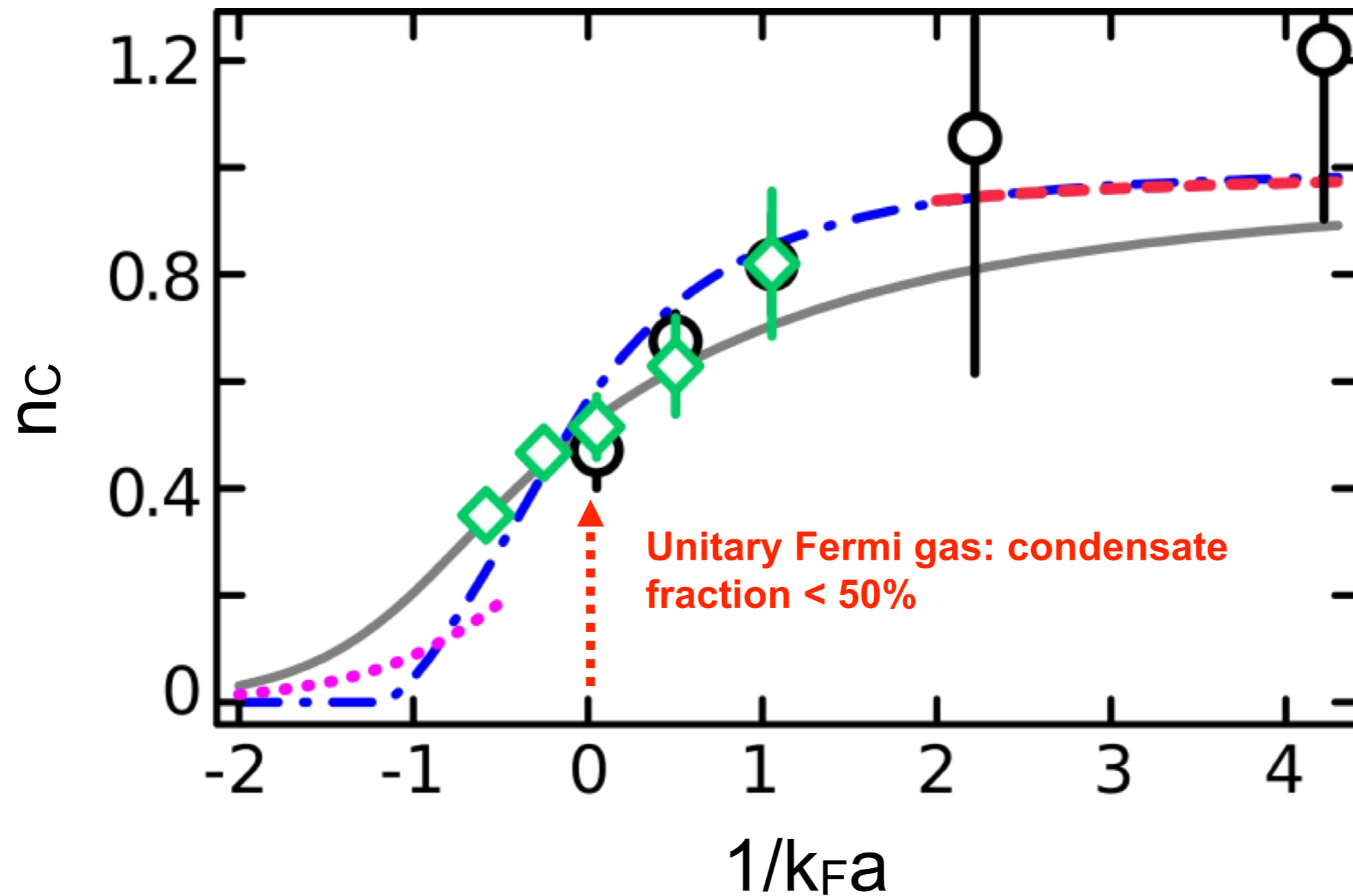


$$I = I_c \sin \varphi - G \Delta \mu$$

Current (μA)



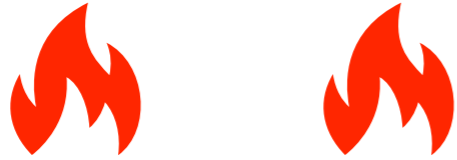
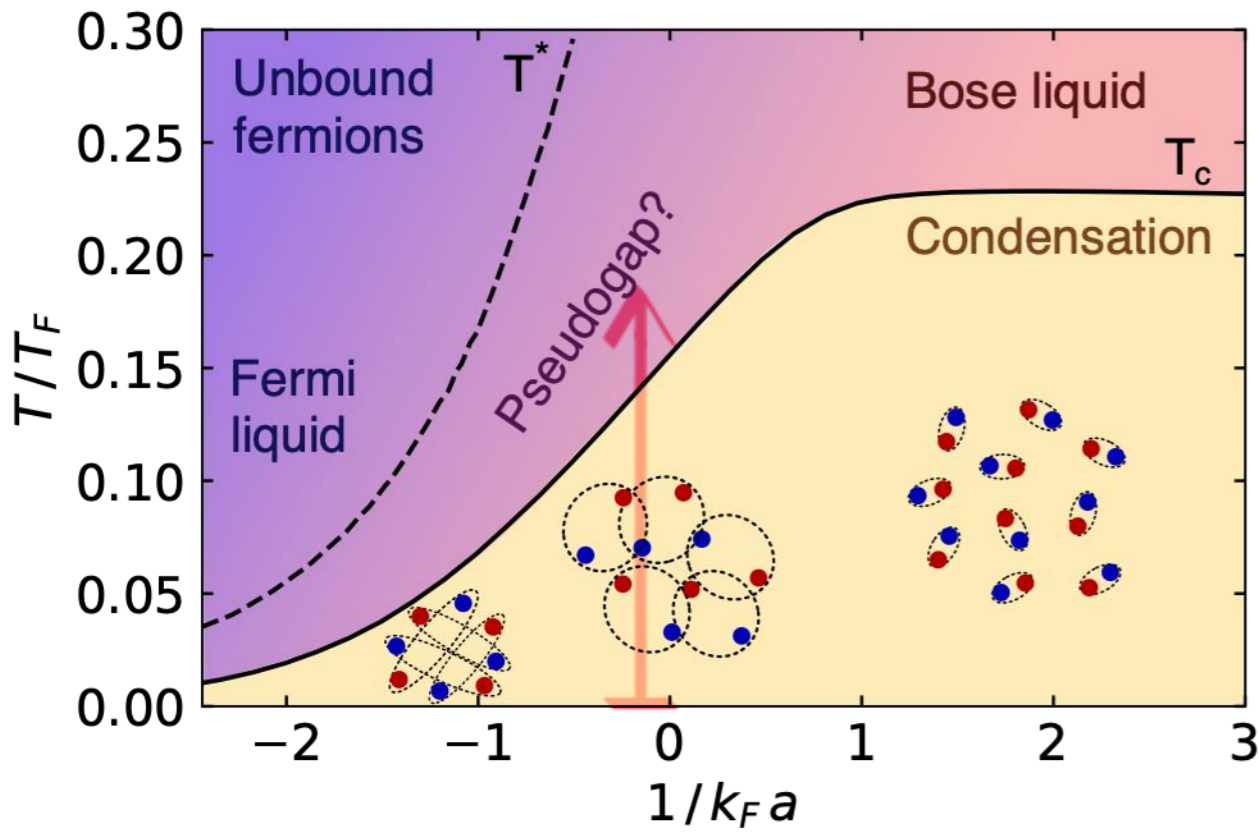




Unitary Fermi gas: condensate fraction < 50%

- \circ average of all data with $V_0/\mu > 0.6$
- \diamond $V_0/E_F \approx 1.06$ data set only
- homogeneous Luttinger-Ward [1]
- - - Quantum Monte Carlo simulations [2]
- BCS approximation [3]
- . - Bogoliubov approximation [4]

- [1] Haussmann et al., *PRA* (2007).
- [2] Astrakharchik et al., *PRL* (2005).
- [3] Gorkov, T. Melik-Barkhudarov, *Sov. Phys. JETP* (1961).
- [4] Giorgini, L. P. Pitaevskii, S. Stringari, *RMP* (2008).

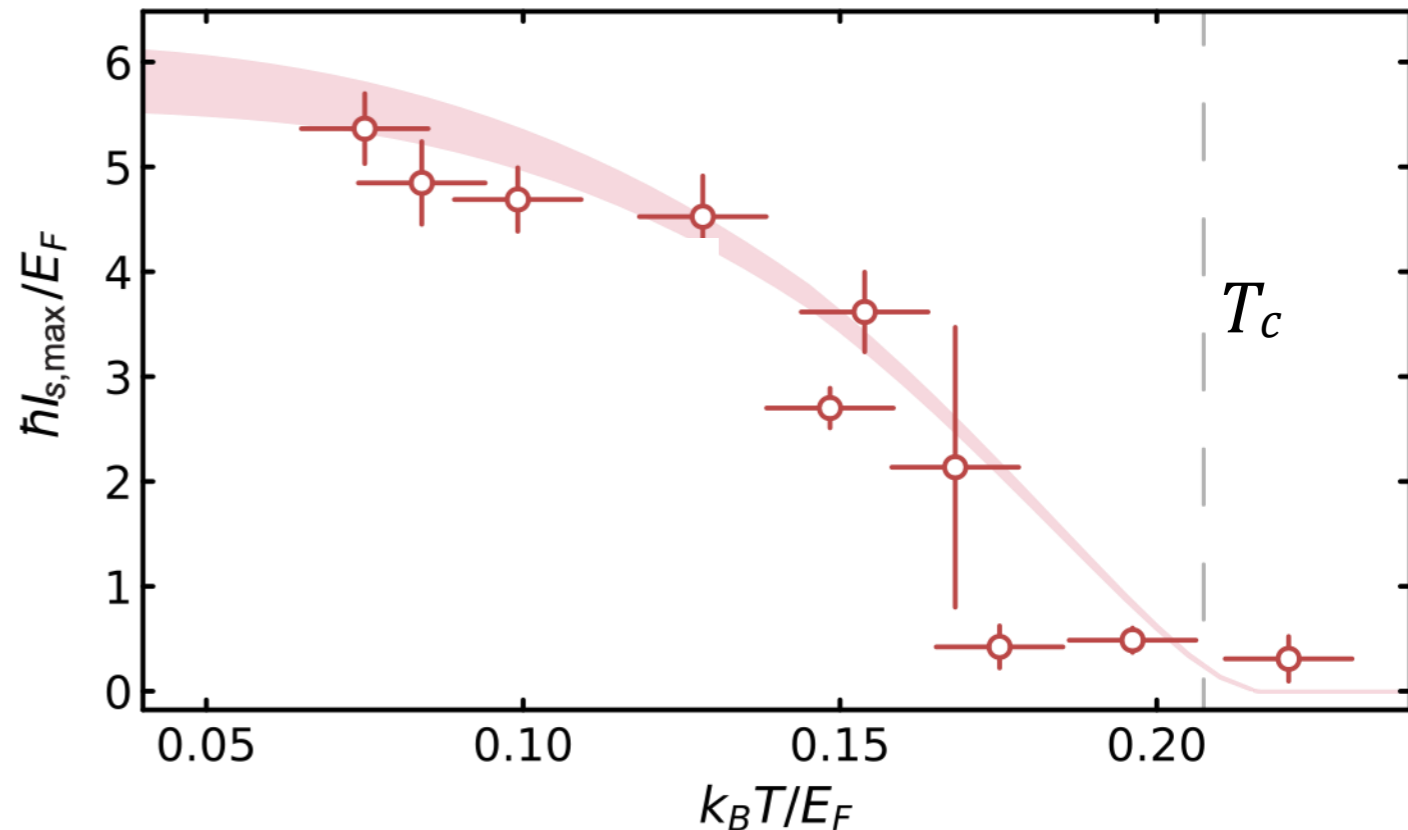
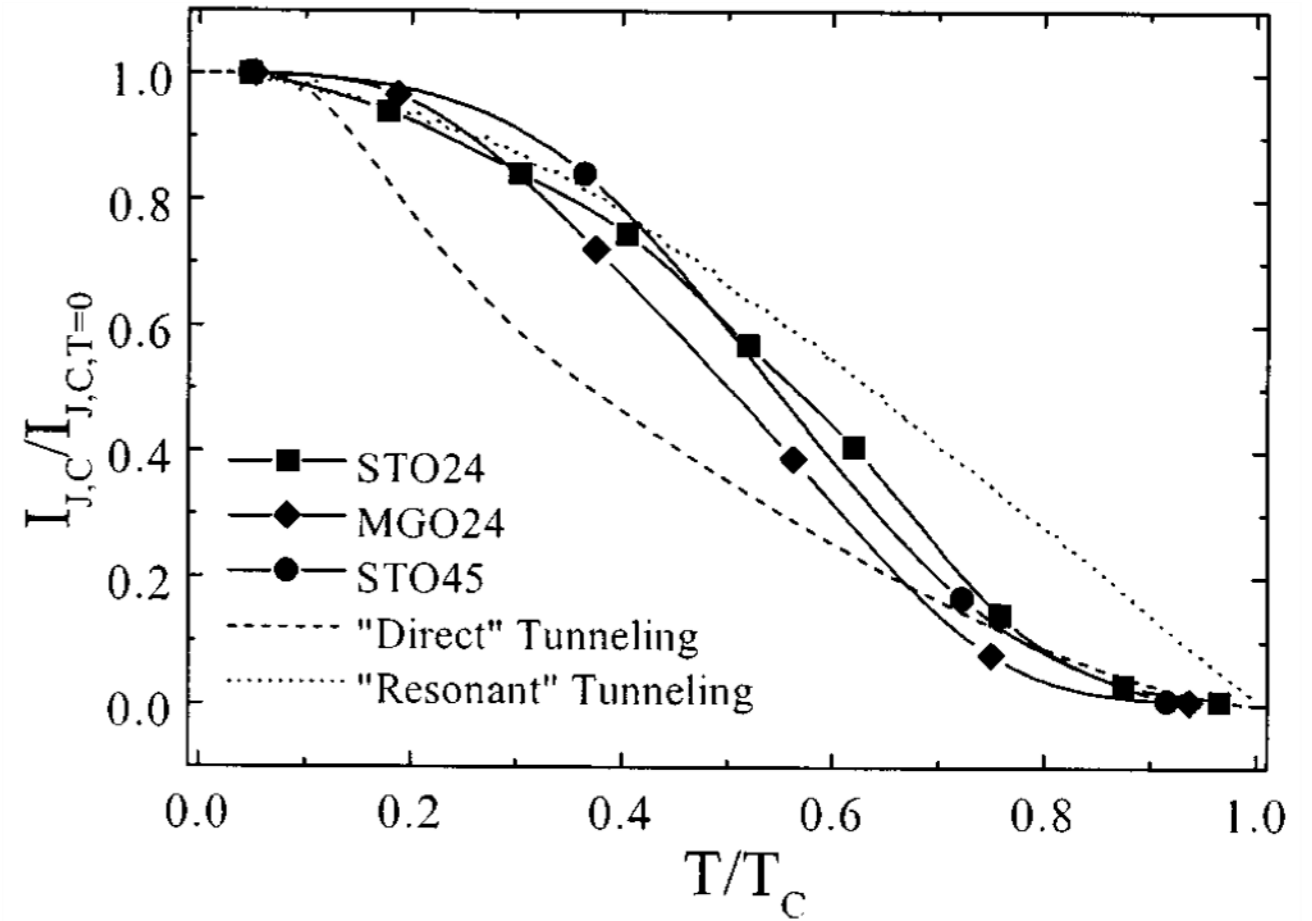


Controlled (parametric) heating

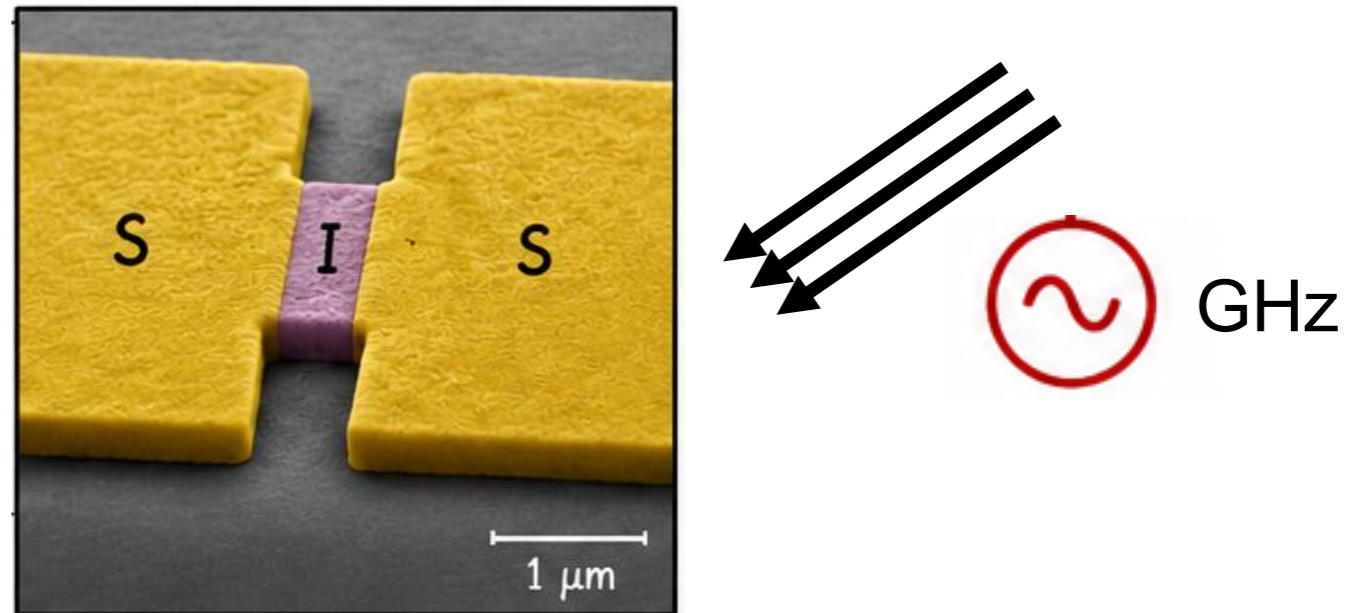
Signature of phase transition in unitary Fermi gas.

Hausmann et al., PRA (2008)

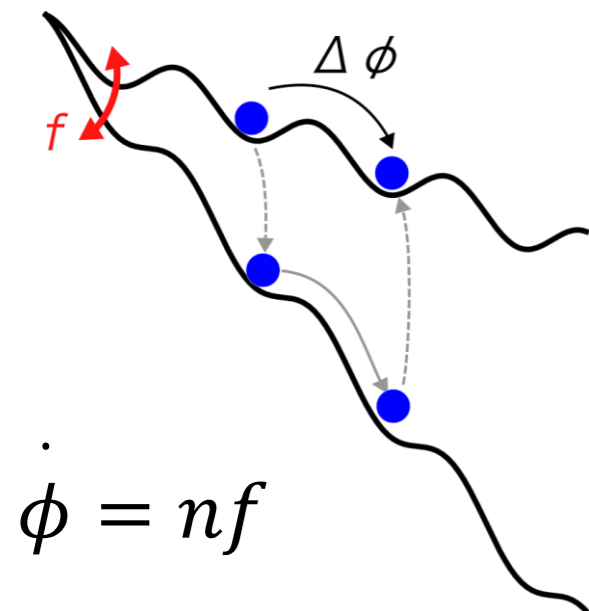
Pini et al., NJP (2020)



AC-drive: Shapiro steps



Irradiating the junction with an oscillatory microwave field: Shapiro steps



$$U(\phi) = -I_{ext}(t)\phi + I_c \cos(\phi)$$

$$I_{ext}(t) = I_0 + I_1 \cos(2\pi f t)$$

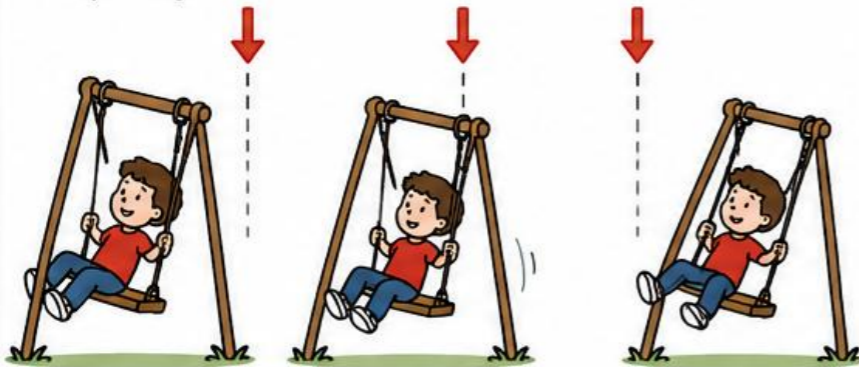
The phase particle jumps by n minima over one modulation period, when its velocity $\dot{\phi}$ is **resonant** with a multiple of the driving frequency f . The phase jumps by $2\pi n$ every modulation period.

AC-drive: Shapiro steps

SHAPIRO STEPS: THE SWING ANALOGY

1. NO SYNCHRONIZATION

You push the swing at a different frequency.

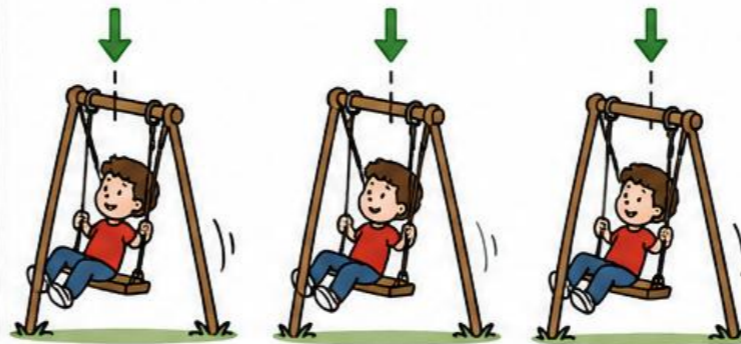


external pushes

The pushes come at random phases.
The swing does not lock to the driver.

2. SYNCHRONIZATION (SHAPIRO STEP)

You push the swing at the right frequency.
Every push arrives at the same point.



external pushes

The swing locks to the pushes.
Its motion (frequency) becomes constant.

ANALOGY → JOSEPHSON JUNCTION

Swing angle ↔ Phase ϕ

Swing speed ↔ Average $d\phi/dt$

Your pushes ↔ Microwave drive (frequency f)

Synchronization ↔ Shapiro step

Locked speed ↔ Locked voltage

Josephson relation:

$$V = \frac{\hbar}{2e} \frac{d\phi}{dt}$$

3. RESULT: PLATEAUS IN VOLTAGE (SHAPIRO STEPS)

In the experiment we increase the bias current I and measure the voltage V .
When synchronization occurs, the voltage locks to discrete values.



At these currents, the voltage stays constant (plateau):

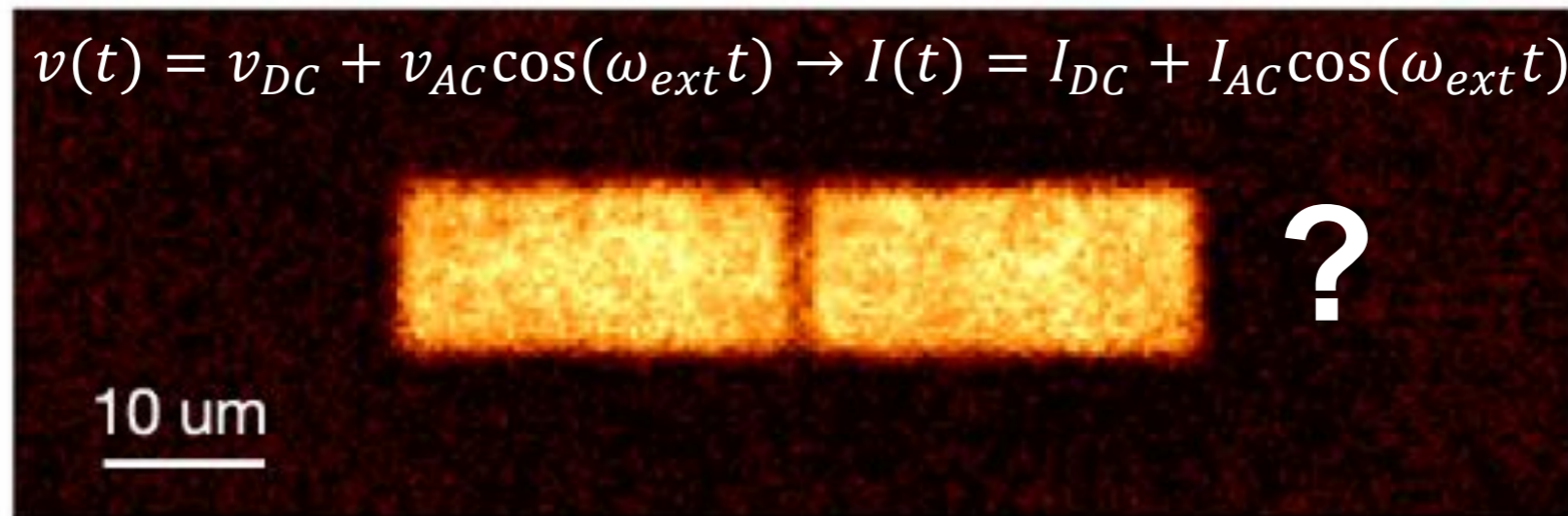
$$V_n = n \frac{hf}{2e}$$

$$n = 1, 2, 3, \dots$$



Shapiro steps appear because the phase (swing) synchronizes with the microwave drive.

Locked frequency → locked voltage.

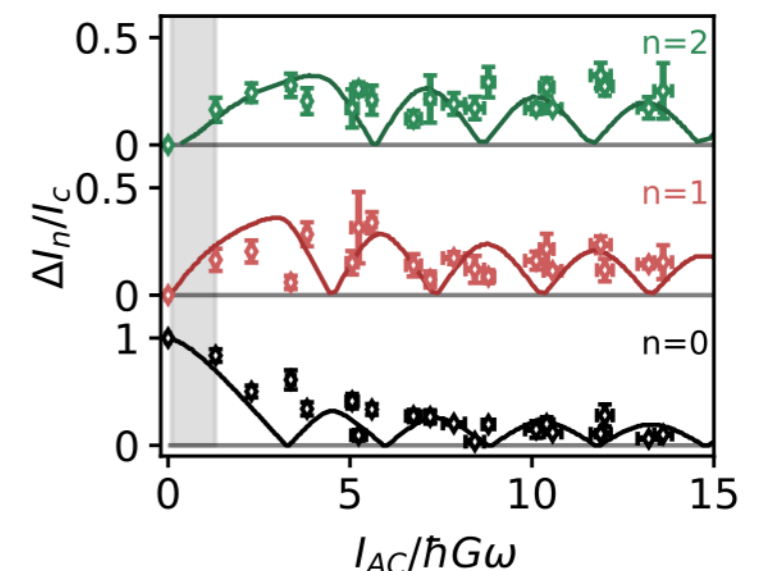
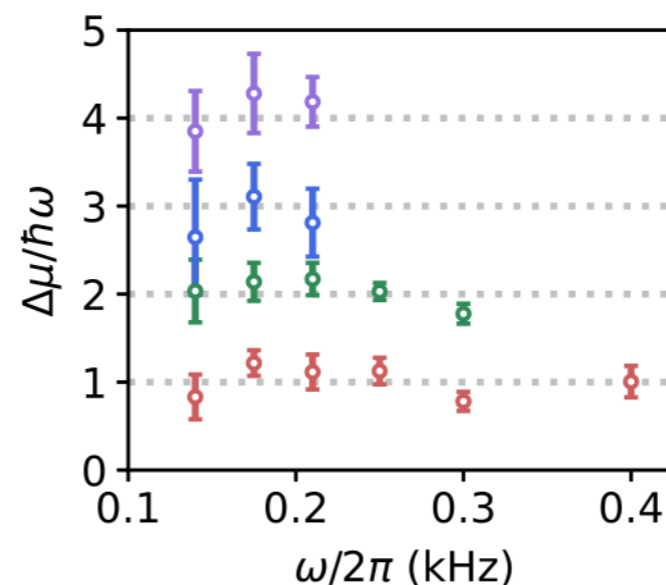
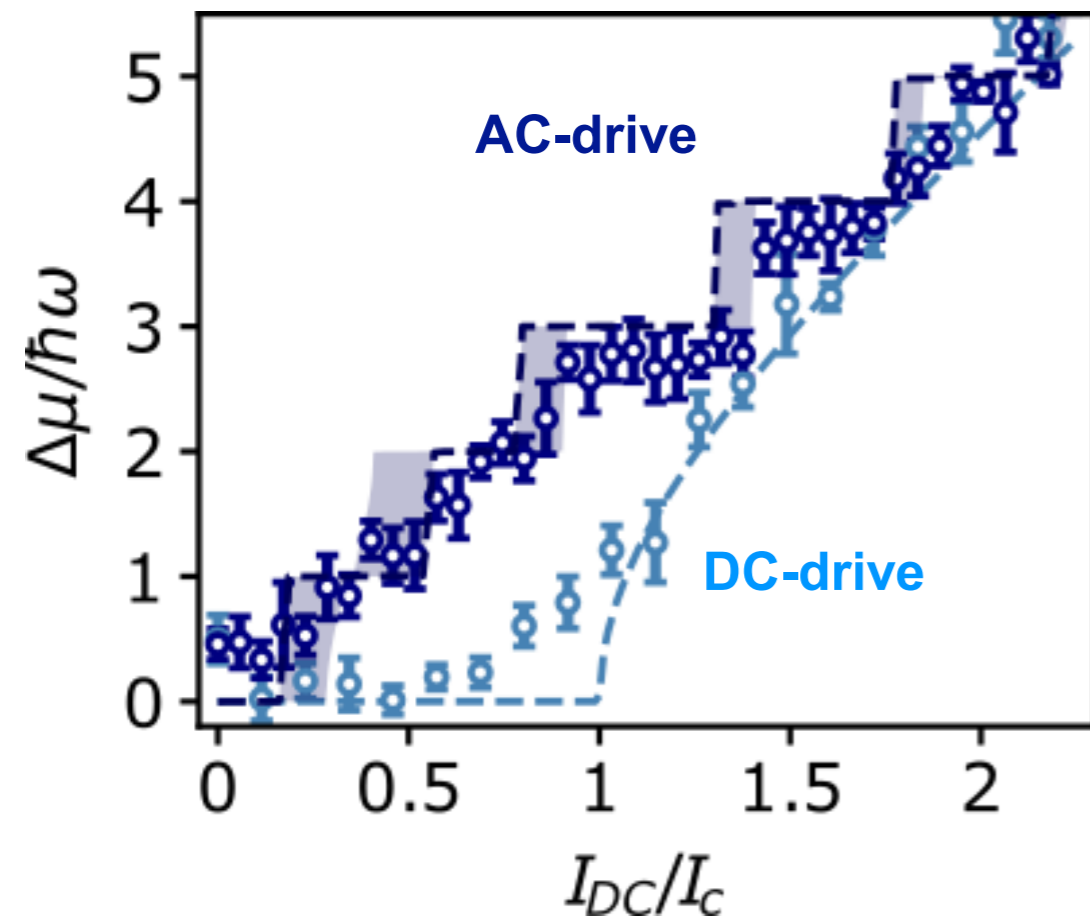


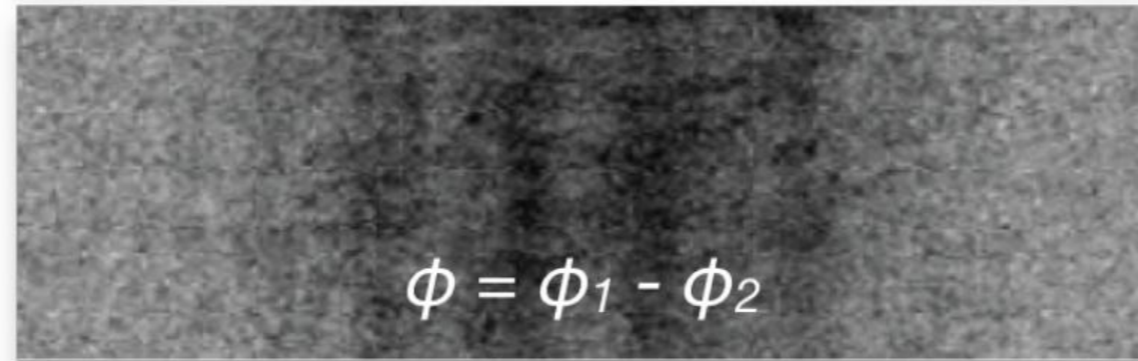
Singh et al, PRL (2024)

Bernhart et al., Science (2025)

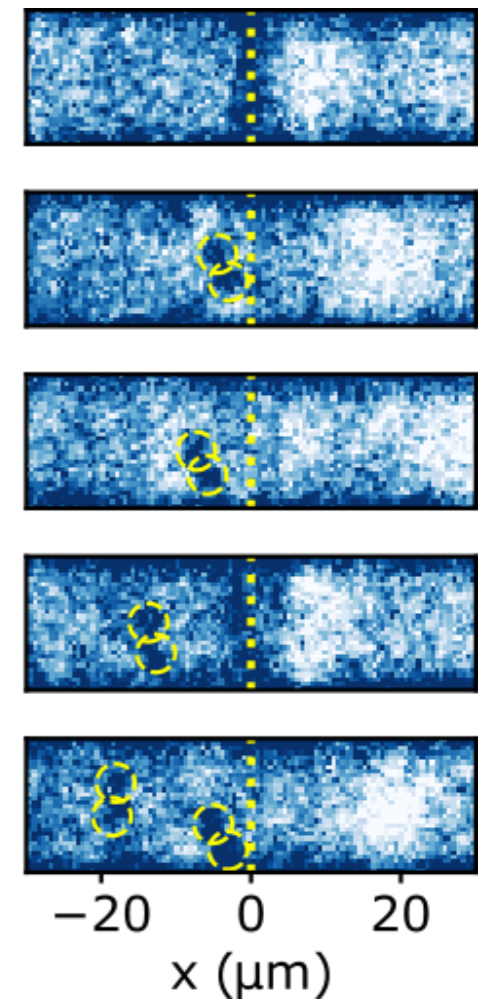
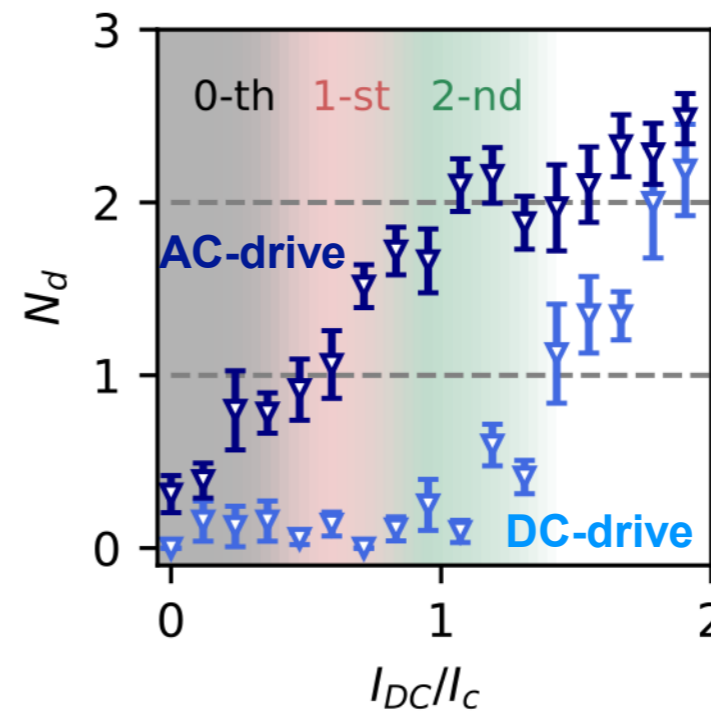
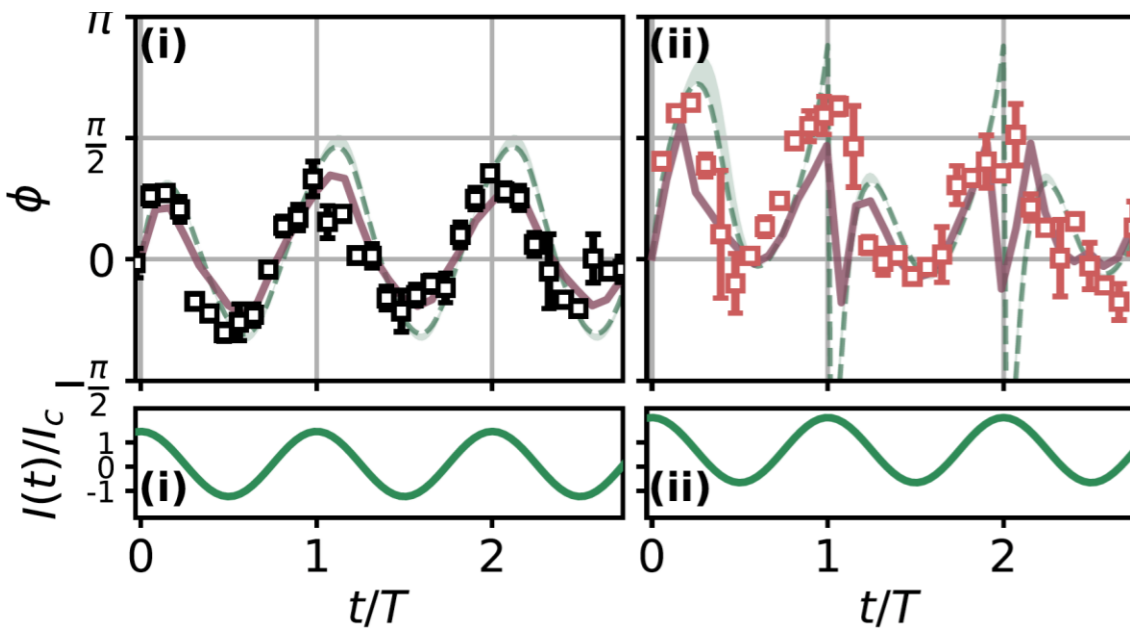
1. Steps in the I-V characteristic

2. Step heights linearly growing with the driving external frequency. Step widths follows Bessel behaviour.





Time-of-flight imaging: relative phase ϕ



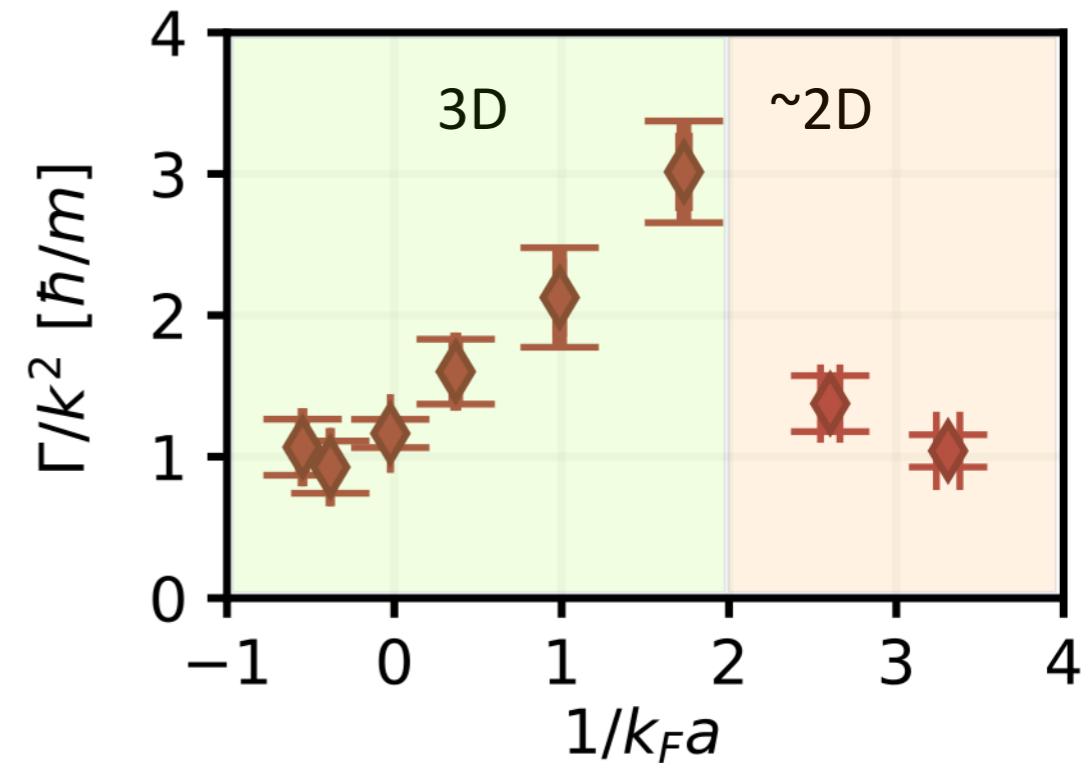
ϕ is synchronised with the external drive. We observe phase-slippage mechanisms as vortex-anti vortex pairs. Impossible for ordinary JJ, due to the different time scale: detection time \gg electrons dynamics time

Other projects:

“*Sagnac*” phonon-interferometer with annular Fermi superfluids (S. Stringari):

Quantum of circulation: angular momentum per particle of BEC-BCS superfluids

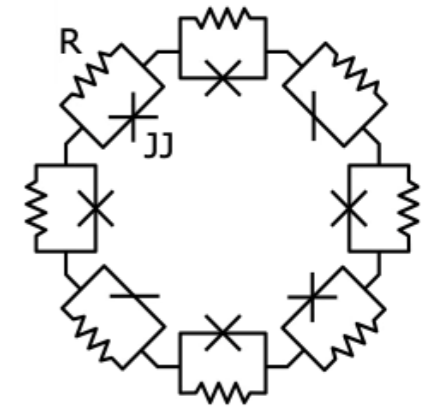
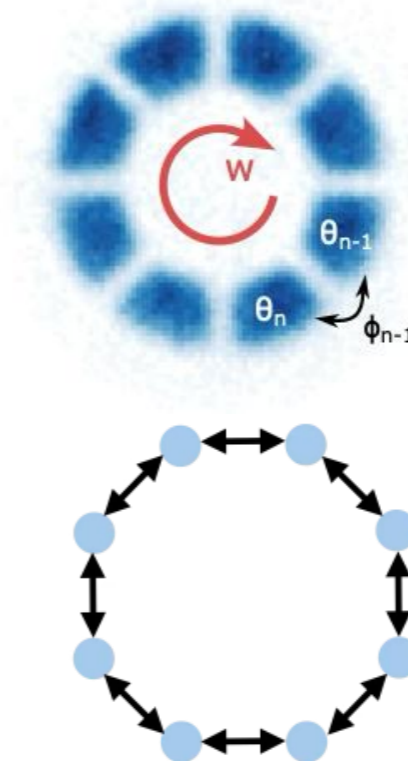
Frometa et al., arXiv:2511.02664v2 (2025), Nat.Phys. in press



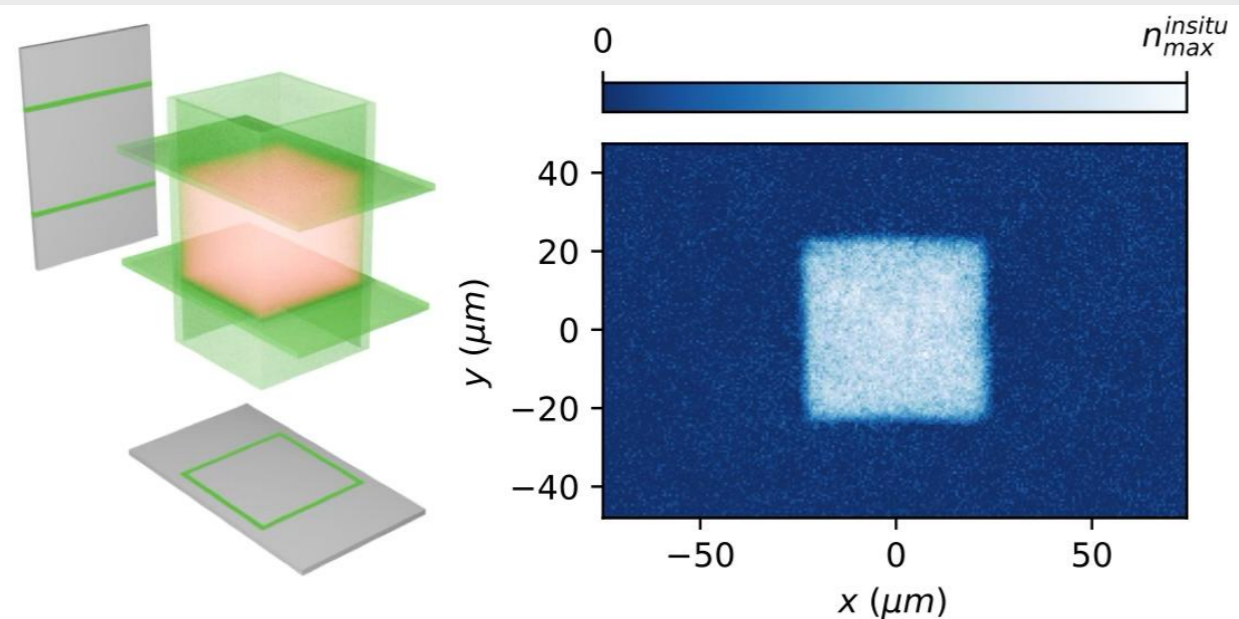
MIT, USTC, Yale, Hamburg, Cambridge, ENS..

Cooperative stabilization of persistent currents in superfluid ring networks: Josephson junction necklace as closed chains of locally coupled **Kuramoto-like oscillators: synchronization phenomena in many-body systems.**

Ciszak et al., arXiv:2601.15121v2 (2026)

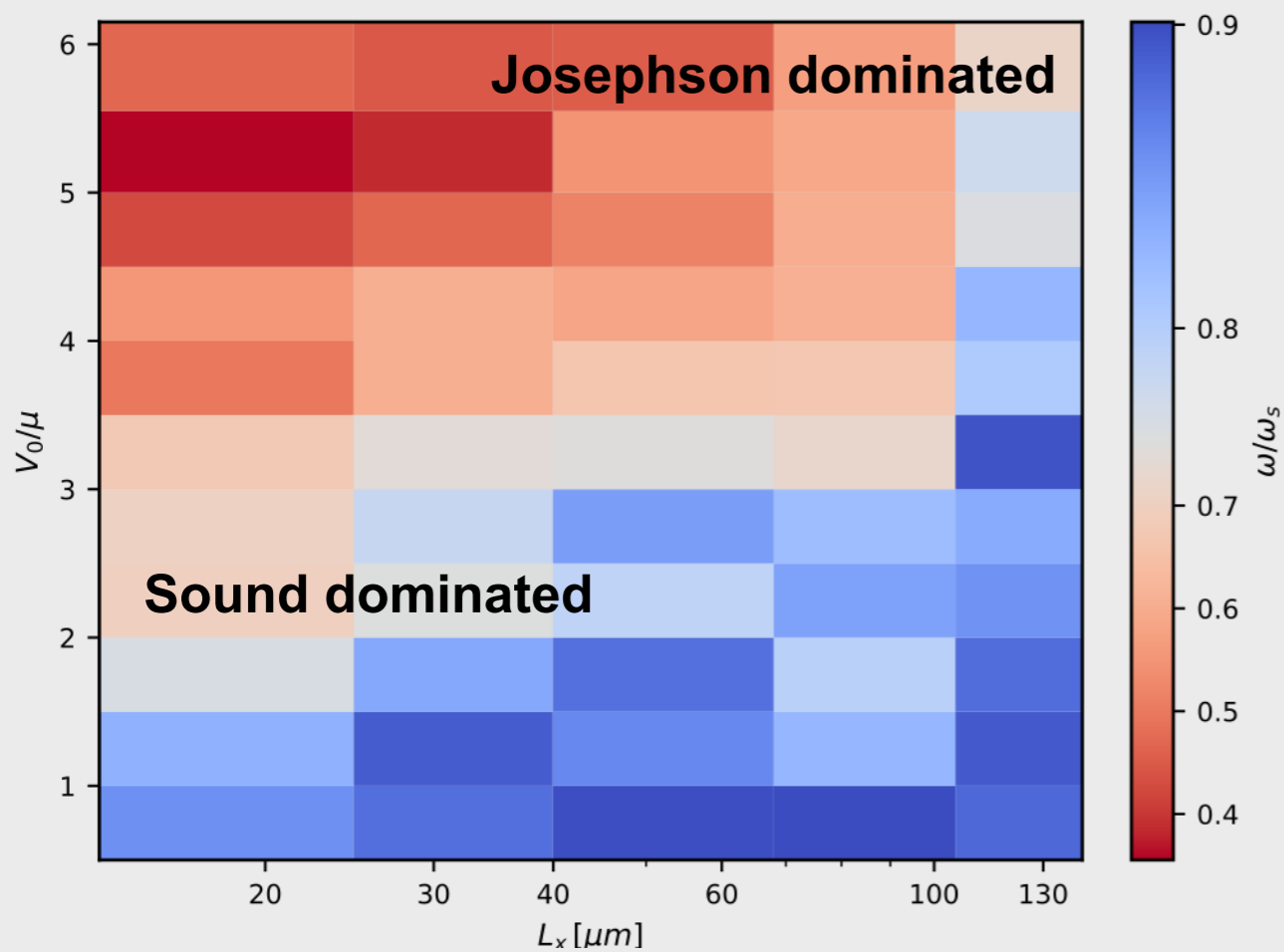


3D Fermi box (with crossed DMDs):

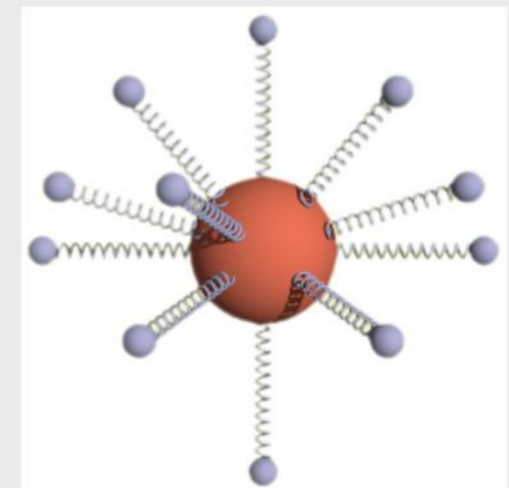
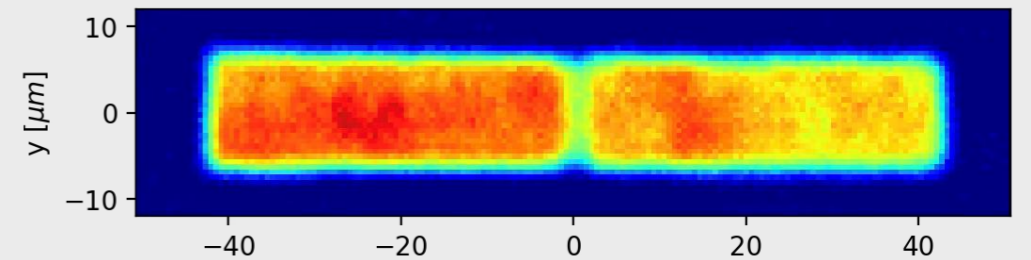


Collaborations and comparison with other Fermi-boxes (MIT, Yale, USTC, Hamburg, EPFL, ...)

Navon et al., Nat. Phys. (2021)



Del Pace et al., *in preparation*



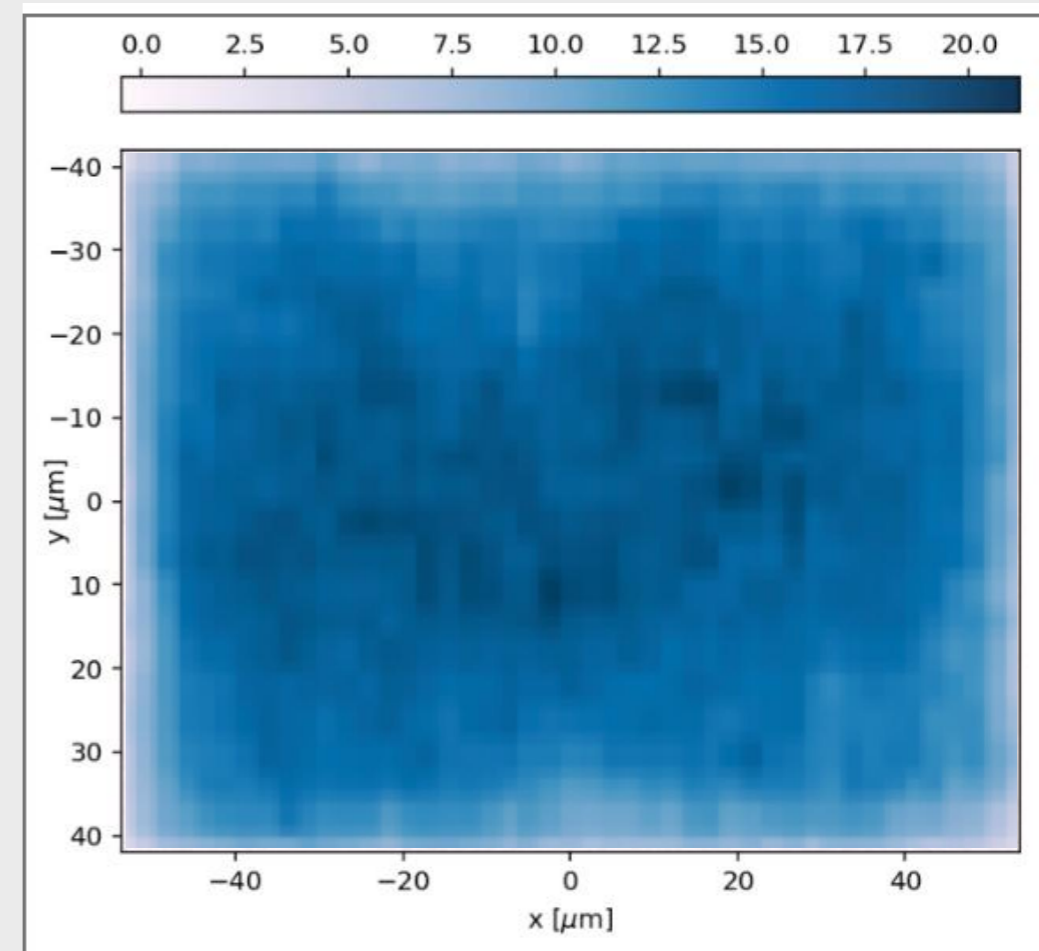
Caldeira and Leggett, Annal. Phys. (1983)

Multiplexed homogenous Fermi boxes:

(Rajkov et al., *in preparation*)

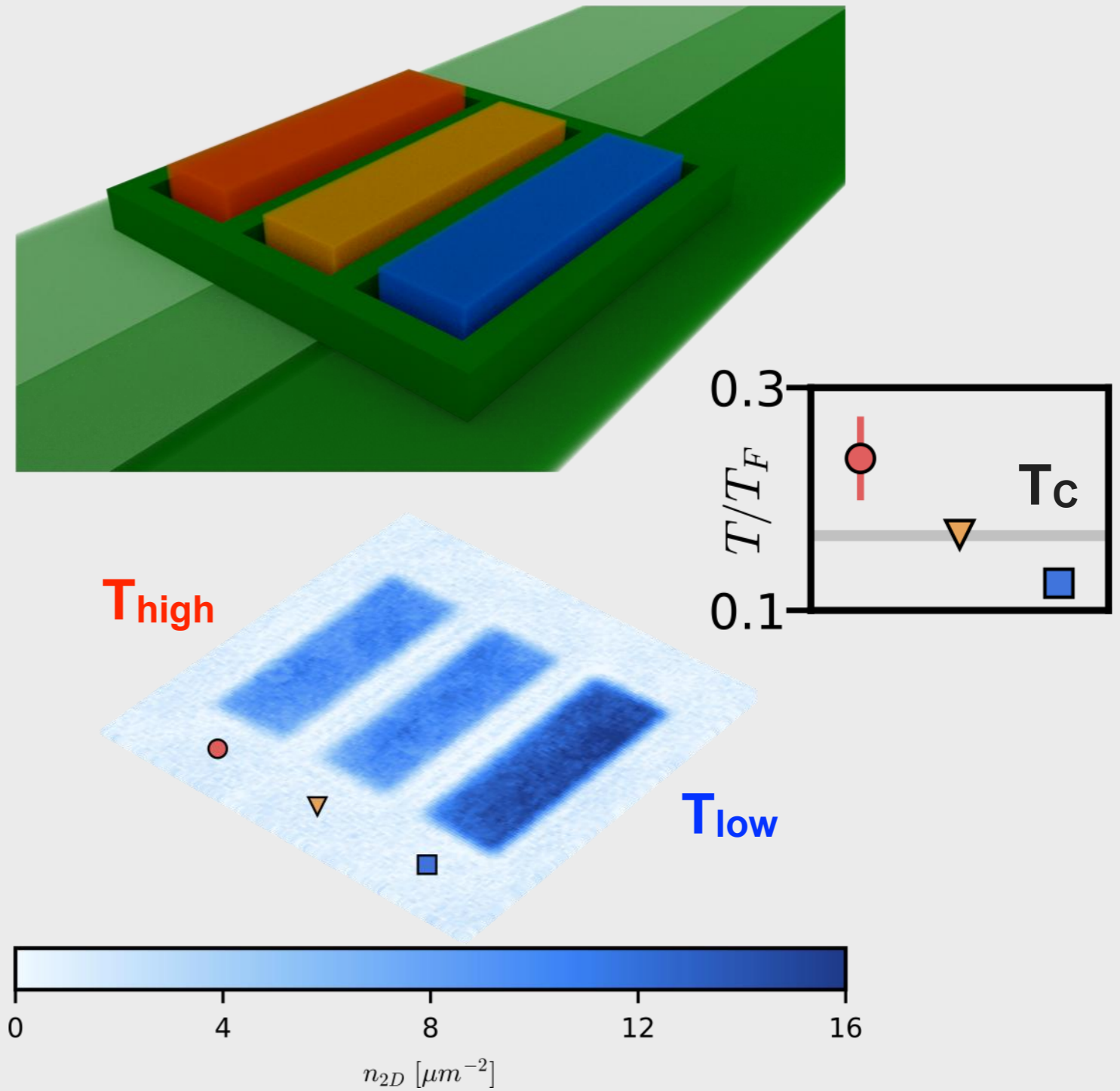
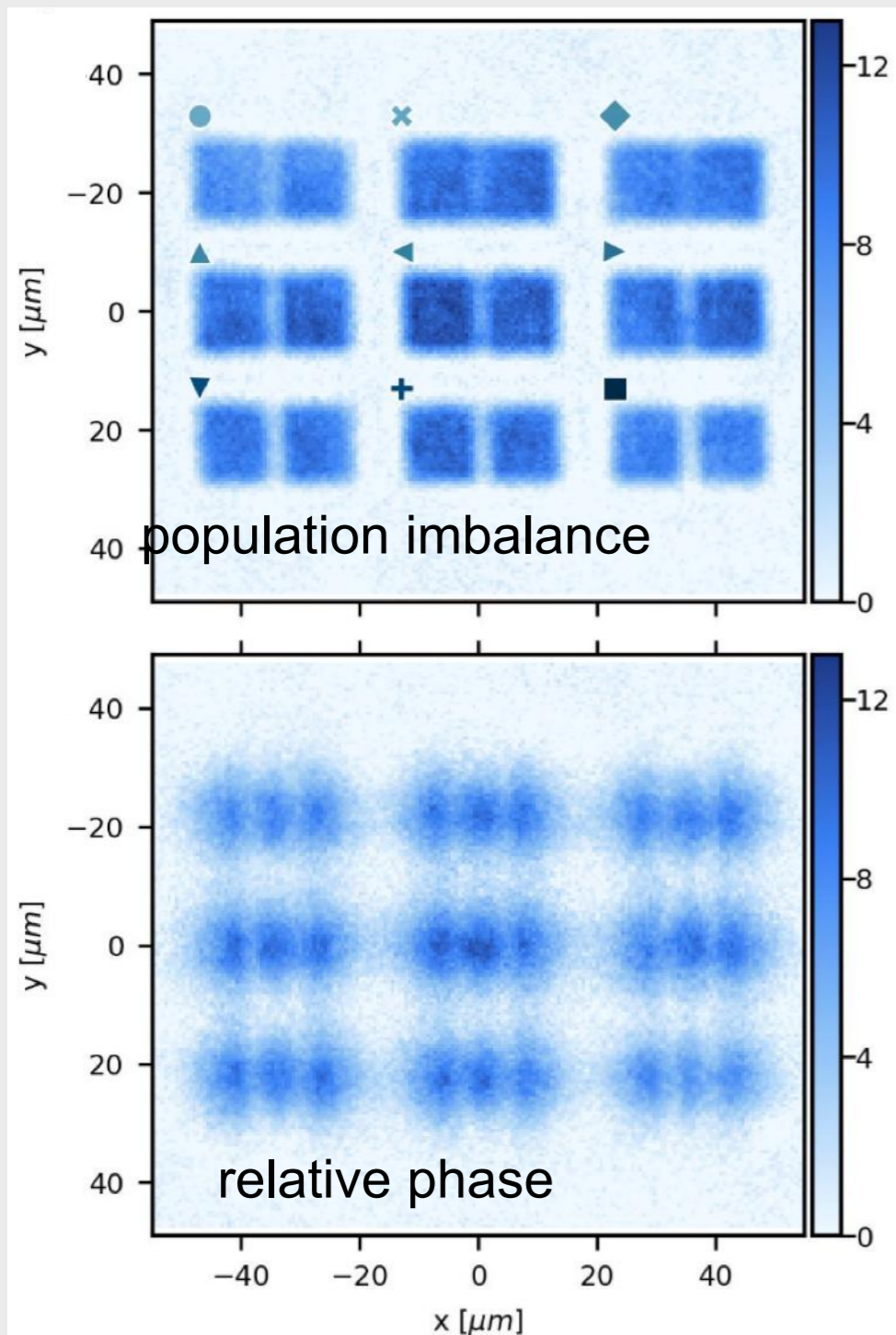
We realise several identical and independent **q-simulation units** with **locally** tunable parameters.

- Parallel execution of the same Hamiltonian with different parameters
- Reconfigurability of the position of each simulating units.



Parallel experiments: ground-state phases, driven dynamics and quantum transport phenomena...

1. Josephson junctions with local phase control.





roati@lens.unifi.it

<http://quantumgases.lens.unifi.it/exp/li>

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